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GLOBAL WEAK SOLUTIONS FOR THE COMPRESSIBLE ACTIVE LIQUID CRYSTAL SYSTEM

GUI-QIANG G. CHEN, APALA MAJUMDAR, DEHUA WANG, AND RONGFANG ZHANG

ABSTRACT. We study the hydrodynamics of compressible flows of active liquid crystals in the Beris-Edwards hydrodynamics framework, using the Landau-de Gennes Q -tensor order parameter to describe liquid crystalline ordering. We prove the existence of global weak solutions for this active system in three space dimensions by the three-level approximations and weak convergence argument. New techniques and estimates are developed to overcome the difficulties caused by the active terms.

1. INTRODUCTION

Nematic liquid crystals are classical examples of complex liquids with long-range orientational order or anisotropic liquids with distinguished directions of average molecular alignment [20, 54]. Nematic order, often described in terms of the collective alignment of constituent elongated particles (*e.g.*, molecules), is ubiquitous; we see collective motion or coordinated motion at all scales ranging from micro-organisms to traffic and flocks of animals. The phrase *active hydrodynamics* is often used to describe the collective dynamics of particles that are constantly maintained out of equilibrium by internal energy sources [26, 38]. Active systems are quite generic in nature, including many biophysical systems such as microtubule bundles [50], dense suspensions of microswimmers [57], bacteria [10], among others. Furthermore, the collective oriented motion is often induced by the elongated shapes of the constituent particles, and hence a large class of active systems are referred to as *active liquid crystals*, especially at high concentrations. We refer to [4, 7, 11, 25, 26, 28, 36, 44, 45, 48] and the references cited therein for more applications and discussions. Active nematics are fundamentally different from the typical passive nematics in the sense that there is no notion of equilibrium; the constituent particles continuously drive the system out of equilibrium leading to striking and novel effects such as the occurrence of giant density fluctuations [38, 40, 47], the spontaneous laminar flow [23, 37, 55], unconventional rheological properties [17, 24, 52], low Reynolds number turbulence [26, 57], and exotic spatial and temporal patterns [8, 21, 38, 39, 49].

Whilst active liquid crystals are popular in the theoretical physics community, a rigorous mathematical description of *active nematics* is relatively new. There are phenomenological models for active liquid crystals in [46], and a common approach is to add phenomenological *active* terms to the hydrodynamic theories for nematic liquid crystals. The mathematics

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of nematic liquid crystals has witnessed a renaissance in recent years, and there are different levels of the mathematical description of passive nematic order: molecular variables describing the orientation and position of each molecule [12, 30], a Oseen-Frank vector field representing the unique direction of preferred molecular alignment [18, 42], and a Landau-de Gennes Q -tensor order parameter that can describe primary and secondary directions of nematic alignment along with variations in the degree of nematic order [2, 3, 31, 34, 35]. More precisely, the Landau-de Gennes Q -tensor order parameter is a d -dimensional symmetric and traceless matrix for the d -dimensional case; the isotropic phase is defined by $Q = 0$. In [9], we analyzed active hydrodynamics in an incompressible Beris-Edwards framework, which is defined by two evolution equations – an evolution equation for the velocity/flow field, and an evolution equation for the Q -tensor with anisotropic stresses from the coupling between flow/order and *active* stresses. We established the existence of global weak solutions in two and three space dimensions for the incompressible *active* Beris-Edwards system. The mathematical machinery in [9] relies on the technical tools in [43] for the incompressible Beris-Edwards system, and we developed new techniques to overcome additional analytical difficulties (compared to [43]) owing to the active stresses.

In this paper, we build on the work in [9] to analyze the following system for compressible flows of active nematic liquid crystals [22, 25] in a bounded domain $\mathcal{O} \subset \mathbb{R}^3$:

$$\begin{cases} \partial_t c + (u \cdot \nabla)c = D_0 \Delta c, \\ \partial_t \rho + \nabla \cdot (\rho u) = 0, \\ \partial_t (\rho u) + \nabla \cdot (\rho u \otimes u) + \nabla P(\rho) - \mu \Delta u - (\nu + \mu) \nabla \operatorname{div} u = \nabla \cdot \tau + \nabla \cdot \sigma, \\ \partial_t Q + (u \cdot \nabla)Q + Q \Omega - \Omega Q = \Gamma H[Q, c], \end{cases} \quad (1.1)$$

where c is the concentration of active particles, ρ is the density of the fluid, $u \in \mathbb{R}^3$ is the flow velocity, the nematic tensor order parameter Q is a traceless and symmetric 3×3 matrix, $P = \kappa \rho^\gamma$ denotes the pressure with adiabatic constant $\gamma > 1$ and constant $\kappa > 0$, $D_0 > 0$ is the diffusion constant, $\mu > 0$ and $\nu > 0$ are the viscosity coefficients, $\Gamma^{-1} > 0$ is the rotational viscosity, and $\Omega = \frac{1}{2}(\nabla u - \nabla u^\top)$ is the antisymmetric part of the strain tensor. Moreover, the tensor

$$H[Q, c] := K \Delta Q - \frac{k}{2}(c - c_*)Q + b(Q^2 - \frac{\operatorname{tr}(Q^2)}{3}I_3) - c_*Q \operatorname{tr}(Q^2)$$

describes the relaxational dynamics of the nematic phase, which can be obtained from the Landau-de Gennes free energy, *i.e.*, $H_{\alpha\beta} = -\frac{\delta \mathcal{F}}{\delta Q_{\alpha\beta}}$ with

$$\mathcal{F} = \int \left(\frac{k}{4}(c - c_*)\operatorname{tr}(Q^2) - \frac{b}{3}\operatorname{tr}(Q^3) + \frac{c_*}{4}|\operatorname{tr}(Q^2)|^2 + \frac{K}{2}|\nabla Q|^2 \right) dA,$$

where K is the elastic constant for the one-constant elastic energy density, c_* is the critical concentration for the isotropic-nematic transition, and $k > 0$ and $b \in \mathbb{R}$ are material-dependent constants. Without loss of generality, we take $K = k = \kappa = 1$ in this paper. The stress tensor $\sigma = (\sigma^{ij})$ has two contributions:

$$\sigma^{ij} = \sigma_r^{ij} + \sigma_a^{ij},$$

with

$$\sigma_r^{ij} = Q^{ik} H^{kj}[Q, c] - H^{ik}[Q, c] Q^{kj},$$

$$\sigma_a^{ij} = \sigma_* c^2 Q^{ij},$$

where σ_r^{ij} is the stress due to the nematic elasticity, and σ_a^{ij} is the active contribution which describes contractile ($\sigma_* > 0$) or extensile ($\sigma_* < 0$) stresses exerted by the active particles along the director field. The symmetric additional stress tensor $\tau = (\tau_{ij})$ is denoted by

$$\tau^{ij} = F(Q)\delta_{ij} - \partial_j Q^{kl} \partial_i Q^{kl} = F(Q)\delta_{ij} - (\nabla Q \odot \nabla Q)^{ij}$$

with

$$F(Q) = \frac{1}{2} |\nabla Q|^2 + \frac{1}{2} \text{tr}(Q^2) + \frac{c_*}{4} \text{tr}^2(Q^2).$$

Here and elsewhere, we use the Einstein summation convention, *i.e.*, we sum over the repeated indices.

We rewrite system (1.1) as

$$\partial_t c + (u \cdot \nabla) c = D_0 \Delta c, \quad (1.2)$$

$$\partial_t \rho + \nabla \cdot (\rho u) = 0, \quad (1.3)$$

$$\begin{aligned} \partial_t (\rho u) + \nabla \cdot (\rho u \otimes u) + \nabla p^\gamma &= \mu \Delta u + (\nu + \mu) \nabla \text{div} u + \nabla \cdot (F(Q) I_3 - \nabla Q \odot \nabla Q) \\ &\quad + \nabla \cdot (Q \Delta Q - \Delta Q Q) + \sigma_* \nabla \cdot (c^2 Q), \end{aligned} \quad (1.4)$$

$$\partial_t Q + (u \cdot \nabla) Q + Q \Omega - \Omega Q = \Gamma H[Q, c], \quad (1.5)$$

with

$$H[Q, c] = \Delta Q - \frac{c - c_*}{2} Q + b \left(Q^2 - \frac{\text{tr}(Q^2)}{3} I_3 \right) - c_* Q \text{tr}(Q^2),$$

and $\Gamma > 0, D_0 > 0, \mu > 0, \nu > 0, c_* > 0, b, \sigma_* \in \mathbb{R}, \gamma > \frac{3}{2}, (x, t) \in \mathcal{O} \times \mathbb{R}^+$; subject to the following initial conditions:

$$(c, \rho, \rho u, Q)|_{t=0} = (c_0, \rho_0, m_0, Q_0)(x) \quad \text{for } x \in \mathcal{O} \subset \mathbb{R}^3, \quad (1.6)$$

with

$$\begin{aligned} c_0 &\in H^1(\mathcal{O}), \quad 0 < \underline{c} \leq c_0 \leq \bar{c} < \infty, \\ Q_0 &\in H^1(\mathcal{O}), \quad Q_0 \in S_0^3 \quad a.e. \text{ in } \mathcal{O}, \end{aligned}$$

and the following boundary conditions on $\partial\mathcal{O}$ with unit outward normal \vec{n} :

$$\nabla c \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \quad u|_{\partial\mathcal{O}} = 0, \quad \nabla Q \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \quad (1.7)$$

satisfying the following compatibility conditions:

$$\rho_0 \in L^\gamma(\mathcal{O}), \quad \rho_0 \geq 0; \quad m_0 \in L^1(\mathcal{O}), \quad m_0 = 0 \text{ if } \rho_0 = 0; \quad \frac{|m_0|^2}{\rho_0} \in L^1(\mathcal{O}). \quad (1.8)$$

The hydrodynamic equations in (1.1), or (1.2)–(1.5), are from [26] with some differences, primarily for technical reasons. Namely, in the concentration equation, the diffusion constants are assumed to be the same in all directions, and the active current is assumed to be zero, which is equivalent to setting $\alpha_1 = 0$ in equations (15a)–(15c) in [26]. Furthermore, the flow-aligning parameter λ in [26] is assumed to be zero; this is not a severe restriction, but just implies that we are in the flow-tumbling regime. Finally, we have also neglected one of the terms in the passive “nematic” stress which does not feature for the two-dimensional systems but can play a role for the three-dimensional systems. Despite these simplifications compared to the successful model presented in [26], our work is a first

step in the rigorous analysis of initial-boundary value problems for compressible active nematics in two and three space dimensions, and the mathematical approach developed here is different from the previous approaches in [9] and [43].

In the simplified system above, the fluid flow is not just dictated by the compressible Navier-Stokes equations, the particle concentration in the fluid and the evolution of the order parameter Q are governed by the parabolic-type equations, with extra nonlinear coupling terms as forcing terms. The term, $F(Q)$, is added to close the energy in our compressible system. Since our system reduces to the compressible Navier-Stokes system in the absence of the concentration c and the Q -tensor, the best result we could expect can not be better than those in [14–16], in which the existence of finite-energy weak solutions of the compressible Navier-Stokes system (allowing initial vacuum) was proved for $\gamma > \frac{3}{2}$. In this paper, our aim is to prove the existence of global weak solutions of this compressible coupled system (1.2)–(1.8) in three space dimensions. In our system, owing to the varying concentration $c = c(x, t)$, we multiply the Q -tensor equation by $-(\Delta Q - Q - c_* Q \operatorname{tr}(Q^2))$, rather than $-H[Q, c]$, to avoid dealing with the interaction terms of the concentration and the Q -tensor and obtain the dissipation and *a priori* estimates for the system. Moreover, the cubic term of the Q -tensor does not appear in the energy with this strategy, so that a positive total energy for this system can be obtained, unlike in [9] where a specific positive energy for the system is re-defined by using the property of the Q -tensor. Furthermore, the highly nonlinear terms in this system cause new mathematical difficulties compared to [9]. However, since the maximum principle holds for the concentration equation (1.2) for c (*i.e.*, c is bounded if the initial condition (3.5) is satisfied; see Lemma 3.2), the highly nonlinear terms can be dealt with via using some cancellation rules as in Lemma A.1 and (2.6) in Proposition 2.1. We remark that the symmetry and tracelessness of the Q -tensor play a key role in the cancellations which are crucial for the proof of the existence of weak solutions. For example, in order to obtain the essential compactness results, the force term in the compressible Navier-Stokes equations should belong to $H^{-1}(\mathcal{O})$ due to Lions [32]. However, the regularity of Q obtained from the Q -tensor equation is $L_t^\infty H_x^1 \cap L_t^2 H_x^2$, which is not enough to achieve this condition. Owing to the cancellations, all the higher order nonlinear terms together vanish, so we do not need to deal with them.

In this paper, we apply the Faedo-Galerkin's method [53] with three levels of approximations to prove the existence of the solutions of the initial-boundary value problem (1.2)–(1.8) in a bounded domain $\mathcal{O} \subset \mathbb{R}^3$. The first level of approximation concerns the artificial pressure due to the possibility of vanishing density and lower integrability of the density. Here we lift the density above zero to avoid the vacuum and add the artificial pressure to increase the integrability of the density. The second level approximation corresponds to the artificial viscosity, which changes the continuity equation from the hyperbolic to parabolic type which ensures higher regularity. The last level is the approximation from the finite-dimensional to infinite-dimensional space. By the weak convergence argument, we obtain the global existence of finite-energy weak solutions defined as follows:

Definition 1.1. For any $T > 0$, (c, ρ, u, Q) is a finite-energy weak solution of problem (1.2)–(1.8) if the following conditions are satisfied:

- (i) $c > 0$, $c \in L^\infty(0, T; L^2(\mathcal{O})) \cap L^2(0, T; H^1(\mathcal{O}))$; $\rho \geq 0$, $\rho \in L^\infty(0, T; L^\gamma(\mathcal{O}))$; $u \in L^2(0, T; H_0^1(\mathcal{O}))$, $Q \in L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}))$, and $Q \in S_0^3$ *a.e.* in $\mathcal{O}_T = [0, T] \times \mathcal{O}$.

- (ii) Equations (1.2)–(1.5) are valid in $\mathcal{D}'(\mathcal{O}_T)$. Moreover, (1.3) is valid in $\mathcal{D}'(0, T; \mathbb{R}^3)$, if (ρ, u) are extended to be zero on $\mathbb{R}^3 \setminus \mathcal{O}$.
- (iii) Energy $E(t)$ is locally integrable on $(0, T)$ and satisfies the energy inequality:

$$\begin{aligned} \frac{d}{dt}E(t) + \frac{D_0}{2}\|\nabla c\|_{L^2}^2 + \frac{\mu}{2}\|\nabla u\|_{L^2}^2 + (\nu + \mu)\|\operatorname{div} u\|_{L^2}^2 + \frac{\Gamma}{2}\|\Delta Q\|_{L^2}^2 + \frac{c_*^2\Gamma}{2}\|Q\|_{L^6}^6 \\ \leq C(\|u\|_{L^2}^2 + \|\nabla Q\|_{L^2}^2 + \|Q\|_{L^2}^2 + \|Q\|_{L^4}^4) \quad \text{in } \mathcal{D}'(0, T), \end{aligned}$$

where

$$E(t) := \int_{\mathcal{O}} \left(\frac{1}{2}|c|^2 + \frac{1}{2}\rho|u|^2 + \frac{\rho^\gamma}{\gamma-1} + \frac{1}{2}|Q|^2 + \frac{1}{2}|\nabla Q|^2 + \frac{c_*}{4}|Q|^4 \right) dx.$$

- (iv) Equation (1.3) is satisfied in the sense of renormalized solutions; that is, for any function $g \in C^1(\mathbb{R})$ with the property:

$$g'(z) \equiv 0 \quad \text{for all } z \geq M \text{ for a sufficiently large constant } M, \quad (1.9)$$

then

$$\partial_t g(\rho) + \operatorname{div}(g(\rho)u) + (g'(\rho)\rho - g(\rho))\operatorname{div} u = 0 \quad \text{in } \mathcal{D}'(0, T). \quad (1.10)$$

Our main result reads as follows.

Theorem 1.1. *Let $\gamma > \frac{3}{2}$ and $\mathcal{O} \subset \mathbb{R}^3$ be a bounded domain of the class $C^{2+\tau}$ for some $\tau > 0$. Assume that the initial data function $(c_0, \rho_0, m_0, Q_0)(x)$ satisfies the compatibility conditions (1.8). Then, for any $T > 0$, problem (1.2)–(1.7) admits a finite-energy weak solution $(c, \rho, u, Q)(t, x)$ on \mathcal{O}_T .*

We shall prove Theorem 1.1 by the aforementioned three-level approximations, including the Faedo-Galerkin approximation, artificial viscosity, artificial pressure, as well as the weak convergence argument, in the spirit of [15, 16]. This approach was used to construct weak solutions to the compressible Beris-Edwards in [56]. As discussed above, new techniques are needed to overcome the difficulties arising from the concentration equation and its coupling with both the fluid and Q -tensor equations. Firstly, by using the Faedo-Galerkin approximation, for any fixed u_n in the finite-dimensional space $C(0, T; X_n)$ (see (3.17)), we obtain a unique solution $(\rho[u_n], c[u_n], Q[u_n])$ of the initial-boundary value problem (3.1)–(3.2) and (3.4). In Lemma 3.2, system (3.13) has complicated interaction terms which cause difficulties in the proof of both the uniqueness of the solution $(c[u_n], Q[u_n])$ and the traceless property of $Q[u_n]$. The tracelessness of the Q -tensor is an important property that guarantees the validity of the cancellation rules in order to treat the highly nonlinear terms caused by the appearance of the concentration and the Q -tensor. In the proof of the tracelessness, we make the L^2 -estimate of $\operatorname{tr} Q$ from an energy inequality, which implies the tracelessness of Q when combined with Gronwall's inequality and the tracelessness of the initial condition of Q . As we mentioned before, the highly nonlinear interaction terms cause difficulties in this procedure. Here we can see that the concentration equation has the maximum principle, which provides the L^∞ -bound for the concentration. Moreover, we show that the solutions to system (3.13) are in $L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}))$, which provides sufficient regularity for the interaction terms so that they stay bounded, and we can prove the uniqueness of the solution and the traceless property of $Q[u_n]$. Then, substituting $(\rho[u_n], c[u_n], Q[u_n])$ into the variational problem of the momentum equation,

we can construct a contraction map and obtain a local solution (ρ_n, c_n, u_n, Q_n) of the approximation system (3.1)–(3.4) on the time interval $[0, T_n]$. Moreover, the global existence of solutions follows from the uniform energy estimate of the approximation system. In order to pass to the limit for solutions (ρ_n, c_n, u_n, Q_n) as $n \rightarrow \infty$ to obtain a solution $(\rho_{\varepsilon, \delta}, c_{\varepsilon, \delta}, u_{\varepsilon, \delta}, Q_{\varepsilon, \delta})$ for the approximation system in the infinite space, we need enough integrability of the solutions. First, the maximum principle satisfied by the concentration equation provides sufficient integrability for the concentration c_n . This, together with the regularity of (u_n, Q_n) obtained in the uniform energy estimate, gives enough compactness for the nonlinear interaction terms of c_n and Q_n in our system. It actually also plays a crucial role when the artificial viscosity and the artificial pressure tend to zero. With the above results and the artificial pressure and viscosity in the approximation system, we have enough regularity and integrability of the density. These integrability and compactness results allow us to pass to the limit as $n \rightarrow \infty$ and obtain a solution for the approximation system in the infinite dimensional space. Secondly, we let the artificial viscosity ε tend to zero to recover the original continuity equation. Here we employ the convergence of the effective viscous flux sequence to deal with the lack of regularity of the density sequence to retrieve the compactness results of the solutions. Lastly, we pass to the limit of the vanishing artificial pressure sequence to obtain a finite-energy weak solution of the original problem, including the vacuum case. Again, we do not have enough integrability for the density. Similarly to the vanishing artificial viscosity procedure, the convergence of the effective viscous flux sequence gives us the much needed higher regularity for the density. We remark that, although the weak convergence argument is similar to that in [15], owing to the extra terms and difficulties, we will provide most of the details for the sake of completeness and convenience to the reader.

The rest of the paper is organized as follows. In §2, we derive the dissipation principle and *a priori* estimates. In §3, we employ the Faedo-Galerkin approximation to obtain a solution of the approximation problem (3.1)–(3.4). In §4, we let the artificial viscosity $\varepsilon \rightarrow 0$ to recover the solution to the hyperbolic continuity equation. In §5, we pass the limit $\delta \rightarrow 0$ in the artificial pressure to complete the proof of Theorem 1.1. In Appendix A, we list some important lemmas and state some important preliminaries that we use intensively in the paper.

2. THE DISSIPATION PRINCIPLE AND A PRIORI ESTIMATES

In this section, we derive the dissipation principle and *a priori* estimates in Proposition 2.1 for system (1.2)–(1.5).

For the sake of convenience, we first introduce some notations. We denote the Sobolev space by $H^k(\mathcal{O})$ for integer $k \geq 1$, equipped with norm $\|\cdot\|_{H_x^k}$, and $H^{-k}(\mathcal{O})$ is the dual space of $H_0^k(\mathcal{O})$. Denote by (\cdot, \cdot) the inner product in $L^2(\mathcal{O})$: If a and b are vectors, then

$$(a, b) = \int_{\mathcal{O}} a(x) \cdot b(x) \, dx;$$

and if A, B are matrices, then

$$(A, B) = \int_{\mathcal{O}} A : B \, dx = \int_{\mathcal{O}} \text{tr}(AB) \, dx$$

with $A : B = \text{tr}(AB)$. Denote by $S_0^3 \subset \mathbb{M}^{3 \times 3}$ the space of Q -tensors in \mathbb{R}^3 , that is,

$$S_0^3 := \{Q \in \mathbb{M}^{3 \times 3} : Q^{ij} = Q^{ji}, \text{tr}(Q) = 0, i, j = 1, 2, 3\}.$$

Define the norm of a matrix by using the Frobenius norm denoted by

$$|Q|^2 := \text{tr}(Q^2) = Q^{ij}Q^{ij}.$$

With respect to this norm, we define the Sobolev spaces for the Q -tensors:

$$H^1(\mathcal{O}, S_0^3) := \left\{ Q : \mathcal{O} \rightarrow S_0^3 : \int_{\mathcal{O}} (|\nabla Q(x)|^2 + |Q(x)|^2) dx < \infty \right\}.$$

Set $|\nabla Q|^2 := \partial_k Q^{ij} \partial_k Q^{ij}$ and $|\Delta Q|^2 := \Delta Q^{ij} \Delta Q^{ij}$.

Proposition 2.1. *Let (c, ρ, u, Q) be a smooth solution of problem (1.3)–(1.8). Then there exists a constant $C > 0$ depending only on $(D_0, b, c_*, \sigma_*, \mu, \nu, \Gamma)$ and the initial data such that, for a given $T > 0$,*

$$\begin{aligned} \frac{d}{dt} E(t) + \frac{D_0}{2} \|\nabla c\|_{L^2}^2 + \frac{\mu}{2} \|\nabla u\|_{L^2}^2 + (\nu + \mu) \|\text{div } u\|_{L^2}^2 + \frac{\Gamma}{2} \|\Delta Q\|_{L^2}^2 + \frac{c_*^2 \Gamma}{2} \|Q\|_{L^6}^6 \\ \leq C (\|u\|_{L^2}^2 + \|\nabla Q\|_{L^2}^2 + \|Q\|_{L^2}^2 + \|Q\|_{L^4}^4) \quad \text{for all } t \in (0, T). \end{aligned} \quad (2.1)$$

In addition, if $(c_0, \rho_0, m_0, Q_0)(x) \in L^\infty \times L^\gamma \times L^2 \times H^1$, then

$$\|c\|_{L^\infty(0, T; L^\infty(\mathcal{O}))} + \|c\|_{L^\infty(0, T; L^2(\mathcal{O})) \cap L^2(0, T; H^1(\mathcal{O}))} \leq C, \quad (2.2)$$

$$\|\rho\|_{L^\infty(0, T; L^\gamma(\mathcal{O}))} \leq C, \quad (2.3)$$

$$\|\sqrt{\rho}u\|_{L^\infty(0, T; L^2(\mathcal{O}))}^2 + \frac{\mu}{2} \|\nabla u\|_{L^2(0, T; L^2(\mathcal{O}))}^2 + (\nu + \mu) \|\text{div } u\|_{L^2(0, T; L^2(\mathcal{O}))}^2 \leq C, \quad (2.4)$$

$$\|Q\|_{L^\infty(0, T; H^1(\mathcal{O}) \cap L^4(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O})) \cap L^6(0, T; L^6(\mathcal{O}))} \leq C. \quad (2.5)$$

Proof. The L^∞ -bound for the concentration $c(t, x)$ follows from the maximum principle and its initial condition $c_0 \in L^\infty$. Using the continuity equation (1.3) and the boundary equation (1.7), we have

$$\begin{aligned} & ((\rho u)_t, u) - (\rho u \otimes u, \nabla u) \\ &= \int_{\mathcal{O}} \rho_t |u|^2 dx + \int_{\mathcal{O}} \rho \partial_t u^i u^i dx - \int_{\mathcal{O}} \rho u^i u^j \partial_j u^i dx \\ &= \int_{\mathcal{O}} \rho_t |u|^2 dx + \frac{1}{2} \int_{\mathcal{O}} \rho \partial_t |u|^2 dx + \frac{1}{2} \int_{\mathcal{O}} |u|^2 \text{div}(\rho u) dx \\ &= \int_{\mathcal{O}} \rho_t |u|^2 dx + \frac{1}{2} \frac{d}{dt} \int_{\mathcal{O}} \rho |u|^2 dx - \frac{1}{2} \int_{\mathcal{O}} \rho_t |u|^2 dx + \frac{1}{2} \int_{\mathcal{O}} |u|^2 \text{div}(\rho u) dx \\ &= \frac{1}{2} \frac{d}{dt} \int_{\mathcal{O}} \rho |u|^2 dx + \frac{1}{2} \int_{\mathcal{O}} |u|^2 (\rho_t + \text{div}(\rho u)) dx \\ &= \frac{1}{2} \frac{d}{dt} \int_{\mathcal{O}} \rho |u|^2 dx, \end{aligned}$$

and

$$\begin{aligned} (\rho^\gamma, \text{div } u) &= - \int_{\mathcal{O}} \rho^{\gamma-1} (\rho_t + \nabla \rho \cdot u) dx \\ &= - \frac{1}{\gamma} \frac{d}{dt} \int_{\mathcal{O}} \rho^\gamma dx - \frac{1}{\gamma} \int_{\mathcal{O}} \partial_i \rho^\gamma u^i dx = - \frac{1}{\gamma} \frac{d}{dt} \int_{\mathcal{O}} \rho^\gamma dx + \frac{1}{\gamma} \int_{\mathcal{O}} \rho^\gamma \text{div } u dx, \end{aligned}$$

which gives

$$(\rho^\gamma, \operatorname{div} u) = \frac{1}{1-\gamma} \frac{d}{dt} \int_{\mathcal{O}} \rho^\gamma dx.$$

We take the inner product of equation (1.2) with c , equation (1.4) with u , and equation (1.5) with $-(\Delta Q - Q - c_* Q \operatorname{tr}(Q^2))$ respectively, sum them up, and then integrate by parts over \mathcal{O} to obtain

$$\begin{aligned} & \frac{d}{dt} \int_{\mathcal{O}} \left(\frac{1}{2} |(c, \sqrt{\rho} u)|^2 + \frac{\rho^\gamma}{\gamma-1} + \frac{1}{2} |(Q, \nabla Q)|^2 + \frac{c_*}{4} |Q|^4 \right) dx + D_0 \|\nabla c\|_{L^2}^2 \\ & + \mu \|\nabla u\|_{L^2}^2 + (\nu + \mu) \|\operatorname{div} u\|_{L^2}^2 + \Gamma \|\Delta Q\|_{L^2}^2 + \Gamma \|\nabla Q\|_{L^2}^2 + c_* \Gamma \|Q\|_{L^4}^4 + c_*^2 \Gamma \|Q\|_{L^6}^6 \\ & = -(u \cdot \nabla c, c) - (\nabla \cdot (\nabla Q \odot \nabla Q), u) - (F(Q) I_3, \nabla u) + (\nabla \cdot (Q \Delta Q - \Delta Q Q), u) \\ & - \sigma_*(c^2 Q, \nabla u) + (u \cdot \nabla Q, \Delta Q) - (u \cdot \nabla Q, Q + c_* Q |Q|^2) - (\Omega Q - Q \Omega, \Delta Q) \\ & + (\Omega Q - Q \Omega, Q + c_* Q |Q|^2) + \Gamma \left(\frac{c - c_*}{2} Q, \Delta Q - Q - c_* Q \operatorname{tr}(Q^2) \right) \\ & - b \Gamma(Q^2, \Delta Q) + b \Gamma(Q^2, Q + c_* Q \operatorname{tr}(Q^2)) + 2c_* \Gamma(Q |Q|^2, \Delta Q) \\ & = \sum_{i=1}^{13} \mathcal{I}_i. \end{aligned}$$

First, by Lemma A.1, we have

$$\mathcal{I}_4 + \mathcal{I}_8 = 0.$$

By simple calculation,

$$\mathcal{I}_2 + \mathcal{I}_3 + \mathcal{I}_6 + \mathcal{I}_7 = 0, \quad \mathcal{I}_9 = 0,$$

as shown below:

$$\begin{aligned} & \mathcal{I}_2 + \mathcal{I}_6 + \mathcal{I}_7 \\ & = -(\nabla \cdot (\nabla Q \odot \nabla Q), u) + (u \cdot \nabla Q, \Delta Q) - (u \cdot \nabla Q, Q + c_* Q |Q|^2) \\ & = - \int_{\mathcal{O}} \partial_j (\partial_i Q^{kl} \partial_j Q^{kl}) u^i dx + \int_{\mathcal{O}} u^i \partial_i Q^{kl} \Delta Q^{kl} dx - \int_{\mathcal{O}} u^i \partial_i Q^{kl} (Q^{kl} + c_* Q^{kl} |Q|^2) dx \\ & = - \int_{\mathcal{O}} (\partial_i \partial_j Q^{kl} \partial_j Q^{kl} u^i + \partial_i Q^{kl} \partial_j \partial_j Q^{kl} u^i) dx + \int_{\mathcal{O}} u^i \partial_i Q^{kl} \Delta Q^{kl} dx \\ & \quad - \int_{\mathcal{O}} u^i \partial_i \left(\frac{1}{2} |Q|^2 + \frac{c_*}{4} \operatorname{tr}^2(Q^2) \right) dx \\ & = - \int_{\mathcal{O}} \partial_i \partial_j Q^{kl} \partial_j Q^{kl} u^i dx + \int_{\mathcal{O}} \left(\frac{1}{2} |Q|^2 + \frac{c_*}{4} \operatorname{tr}^2(Q^2) \right) \operatorname{div} u dx \\ & = \int_{\mathcal{O}} \left(\frac{1}{2} |\nabla Q|^2 + \frac{1}{2} |Q|^2 + \frac{c_*}{4} \operatorname{tr}^2(Q^2) \right) \operatorname{div} u dx = -\mathcal{I}_3, \end{aligned} \tag{2.6}$$

and

$$\mathcal{I}_9 = (\Omega Q - Q \Omega, Q + c_* Q |Q|^2) = -(\Omega Q + Q \Omega, Q + c_* Q |Q|^2) + 2(\Omega Q, Q + c_* Q |Q|^2) = 0,$$

where we have used the fact that Q is symmetric and Ω is skew-symmetric.

Moreover, by the Young inequality, we have

$$\begin{aligned}
|\mathcal{I}_1| &= |(u \cdot \nabla c, c)| \leq C \|u\|_{L^2} \|\nabla c\|_{L^2} \|c\|_{L^\infty} \leq \frac{D_0}{4} \|\nabla c\|_{L^2}^2 + C \|u\|_{L^2}^2, \\
|\mathcal{I}_5| &= |\sigma_*(c^2 Q, \nabla u)| \leq C \|c\|_{L^\infty}^2 \|\nabla u\|_{L^2} \|Q\|_{L^2} \leq \frac{\mu}{4} \|\nabla u\|_{L^2}^2 + C \|Q\|_{L^2}^2, \\
|\mathcal{I}_{10}| &= |\Gamma| \left(\frac{c - c_*}{2} Q, \Delta Q - Q - c_* Q \operatorname{tr}(Q^2) \right) \leq \frac{\Gamma}{8} \|\Delta Q\|_{L^2}^2 + C \|Q\|_{L^2}^2 + C \|Q\|_{L^4}^4, \\
|\mathcal{I}_{11}| &= |-b\Gamma(Q^2, \Delta Q)| \leq \frac{\Gamma}{8} \|\Delta Q\|_{L^2}^2 + C \|Q\|_{L^4}^4, \\
|\mathcal{I}_{12}| &= |b\Gamma(Q^2, Q + c_* Q \operatorname{tr}(Q^2))| = \left| b\Gamma \int_{\mathcal{O}} (\operatorname{tr}(Q))^3 dx + bc_* \Gamma \int_{\mathcal{O}} (\operatorname{tr}(Q))^3 |Q|^2 dx \right| \\
&\leq \frac{c^2 \Gamma}{2} \|Q\|_{L^6}^6 + C \|Q\|_{L^2}^2 + C \|Q\|_{L^4}^4, \\
\mathcal{I}_{13} &= 2c_* \Gamma(Q|Q|^2, \Delta Q) = 2c_* \Gamma \int_{\mathcal{O}} \partial_{kk} Q^{ij} Q^{ij} \operatorname{tr}(Q^2) dx \\
&= -2c_* \Gamma \int_{\mathcal{O}} \partial_k Q^{ij} \partial_k Q^{ij} \operatorname{tr}(Q^2) dx - 2c_* \Gamma \int_{\mathcal{O}} \partial_k Q^{ij} Q^{ij} \partial_k \operatorname{tr}(Q^2) dx \\
&= -2c_* \Gamma \int_{\mathcal{O}} |\nabla Q|^2 |Q|^2 dx - c_* \Gamma \int_{\mathcal{O}} |\nabla \operatorname{tr}(Q^2)|^2 dx \leq 0.
\end{aligned}$$

Combining all the above estimates, we obtain the desired result. \square

Corollary 2.1. *For any smooth solution (c, ρ, u, Q) of problem (1.2)–(1.8),*

$$Q \in L^{10}(\mathcal{O}_T) \cap L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O})), \quad \nabla Q \in L^{\frac{10}{3}}(\mathcal{O}_T).$$

Proof. From estimate (2.5) and the Gagliardo-Nirenberg inequality in Lemma A.3, we have

$$\begin{aligned}
\|Q\|_{L^{10}(\mathcal{O})} &\leq C \|Q\|_{L^6(\mathcal{O})}^{\frac{4}{5}} \|\Delta Q\|_{L^2(\mathcal{O})}^{\frac{1}{5}} + C \|Q\|_{L^6(\mathcal{O})}, \\
\|\nabla Q\|_{L^{\frac{10}{3}}(\mathcal{O})} &\leq C \|\nabla Q\|_{L^2(\mathcal{O})}^{\frac{2}{5}} \|\Delta Q\|_{L^2(\mathcal{O})}^{\frac{3}{5}} + C \|\nabla Q\|_{L^2(\mathcal{O})}.
\end{aligned}$$

Then the proof is complete. \square

3. THE FAEDO-GALERKIN APPROXIMATION

In this section, we use the classical Faedo-Galerkin method to construct a solution (c_n, ρ_n, u_n, Q_n) of the initial-boundary value problem (3.1)–(3.4) below, which is an approximation system to (1.2)–(1.5). Since this initial-boundary value problem involves coupling terms with the concentration, velocity, and Q-tensor in the force, such as the highly nonlinear terms $H[Q, c]$, $F(Q)$, and $c^2 Q$, and the terms with high order derivative $Q\Delta Q - \Delta Q Q$, we need to handle all these couplings carefully in order to construct the approximation successfully by using the Faedo-Galerkin method. First, we show that there is a unique solution $(\rho[u_n], c[u_n], Q[u_n])$ to the initial-boundary problem (3.1)–(3.2) and (3.4) for any u_n in the finite-dimensional space $C(0, T; X_n)$. Then, substituting $(\rho[u_n], c[u_n], Q[u_n])$ into the variational problem of the momentum equation, we can obtain a local solution (ρ_n, c_n, u_n, Q_n) of the approximation system (3.1)–(3.4) on the time interval $[0, T_n]$ by using the contraction map theorem. Thanks to the cancellation

rule we mentioned in §2 and the fact that the concentration equation obeys the maximum principle, we can overcome the difficulties caused by the complicated coupling terms of the system both in the proof of extending the local solution to a global one by the uniform energy estimates of the system with respect to n and in the proof of the existence of the first level approximation solution as $n \rightarrow \infty$.

For fixed $\delta > 0$ and $\varepsilon > 0$, we solve the following approximation problem:

$$\partial_t c + (u \cdot \nabla)c = D_0 \Delta c, \quad (3.1)$$

$$\partial_t \rho + \nabla \cdot (\rho u) = \varepsilon \Delta \rho, \quad (3.2)$$

$$\begin{aligned} \partial_t(\rho u) + \nabla \cdot (\rho u \otimes u) + \nabla \cdot (\rho^\gamma + \delta \nabla \rho^\beta) + \varepsilon(\nabla \rho \cdot \nabla)u \\ = \mu \Delta u + (\nu + \mu) \nabla \operatorname{div} u + \nabla \cdot (F(Q)I_3 - \nabla Q \odot \nabla Q) \\ + \nabla \cdot (Q \Delta Q - \Delta Q Q) + \sigma_* \nabla \cdot (c^2 Q), \end{aligned} \quad (3.3)$$

$$\partial_t Q + (u \cdot \nabla)Q + Q \Omega - \Omega Q = \Gamma H[Q, c], \quad (3.4)$$

complemented with the modified initial conditions:

$$c|_{t=0} = c_0 \in H^1(\mathcal{O}), \quad 0 < \underline{c} \leq c_0(x) \leq \bar{c}, \quad (3.5)$$

$$\rho|_{t=0} = \rho_0 \in C^3(\bar{\mathcal{O}}), \quad 0 < \underline{\rho} \leq \rho_0(x) \leq \bar{\rho}, \quad (3.6)$$

$$(\rho u)|_{t=0} = m_0(x) \in C^2(\bar{\mathcal{O}}), \quad (3.7)$$

$$Q|_{t=0} = Q_0(x) \in H^1(\mathcal{O}), \quad Q_0 \in S_0^3 \text{ a.e. in } \mathcal{O}, \quad (3.8)$$

and the boundary conditions on $\partial\mathcal{O}$ with unit outward normal \vec{n} :

$$\nabla c \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \quad \nabla \rho \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \quad (3.9)$$

$$u|_{\partial\mathcal{O}} = 0, \quad \frac{\partial Q}{\partial \vec{n}}|_{\partial\mathcal{O}} = 0, \quad (3.10)$$

where \underline{c} , \bar{c} , $\underline{\rho}$, and $\bar{\rho}$ are positive constants.

Remark 3.1. We now remark on the extra terms in the approximation system. The vanishing viscosity $\varepsilon \Delta \rho$ converts equation (1.3) from the hyperbolic to parabolic type and provides the higher regularity of ρ . The term, $\delta \nabla \rho^\beta$, is added to obtain the higher integrability of ρ for some constant $\beta > 0$. The extra term $\varepsilon \nabla \rho \cdot \nabla u$ is needed to cancel some *bad* terms that do not vanish in the energy estimates.

3.1. The Neumann problem for the density and Q -tensor. We first state the following existence results, which can be found in [15].

Lemma 3.1. *Assume that $u \in C([0, T]; C^2(\bar{\mathcal{O}}, \mathbb{R}^3))$ with $u|_{\partial\mathcal{O}} = 0$. Then there exists the following mapping $S = S[u]$:*

$$S : C([0, T]; C^2(\bar{\mathcal{O}}, \mathbb{R}^3)) \rightarrow C([0, T]; C^3(\bar{\mathcal{O}}))$$

such that

- (i) $\rho = S[u]$ is the unique classical solution of (3.2), (3.6)–(3.7), and (3.9),
- (ii) For all $t \geq 0$,

$$\underline{\rho} \exp \left(- \int_0^t \|\operatorname{div} u(s)\|_{L^\infty(\mathcal{O})} ds \right) \leq S[u](t, x) \leq \bar{\rho} \exp \left(\int_0^t \|\operatorname{div} u(s)\|_{L^\infty(\mathcal{O})} ds \right),$$

(iii) For any u_1 and u_2 in the set:

$$\mathcal{N}_N = \{v : v \in C([0, T]; C_0^2(\mathcal{O}, \mathbb{R}^3)), \|v\|_{C([0, T]; C_0^2(\bar{\mathcal{O}}, \mathbb{R}^3))} \leq N\} \quad (3.11)$$

with some suitable constant $N > 0$,

$$\|S[u_1] - S[u_2]\|_{C([0, T]; H^1(\mathcal{O}))} \leq TC(N, T)\|u_1 - u_2\|_{C([0, T]; H_0^1(\mathcal{O}))}. \quad (3.12)$$

Lemma 3.2. For each $u \in C([0, T]; C_0^2(\bar{\mathcal{O}}, \mathbb{R}^3))$ with $u|_{\partial\mathcal{O}} = 0$, there exists a unique solution of the following initial-boundary value problem:

$$\begin{cases} \partial_t c + (u \cdot \nabla)c = D_0 \Delta c, \\ \partial_t Q + (u \cdot \nabla)Q + Q\Omega - \Omega Q = \Gamma H[Q, c], \\ Q|_{t=0} = Q_0(x) \in H^1(\mathcal{O}), \quad \nabla Q \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \\ c|_{t=0} = c_0(x) \in H^1(\mathcal{O}), \quad \nabla c \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \end{cases} \quad (3.13)$$

with $0 < \underline{c} \leq c_0 \leq \bar{c} < \infty$ such that $Q \in L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}))$, $c \in L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}))$, and $0 < \underline{c} \leq c \leq \bar{c} < \infty$. Moreover, the above mapping $u \mapsto (c[u], Q[u])$ is continuous from \mathcal{N}_N to $(L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}))) \times (L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O})))$, and $Q[u] \in S_0^3$ a.e. in \mathcal{O}_T .

Proof. We divide the proof into five steps.

1. The existence of a solution (c, Q) can be achieved by the standard parabolic theory [29]. The boundedness of c is guaranteed by the fact that the maximum principle is valid for the first equation in system (3.13) and the initial condition of c . Next, we show that (c, Q) belongs to $(L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}))) \times (L^\infty(0, T; H^1(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O})))$.

2. Assume that $u \in \mathcal{N}_N$. First, let us take the inner product of the first equation in (3.13) with $c - \Delta c$, and the second equation in (3.13) with $-(\Delta Q - Q - c_* Q \operatorname{tr}(Q^2))$, then sum them up to obtain

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \int_{\mathcal{O}} (|(c, Q, \nabla c, \nabla Q)|^2 + \frac{c_*}{2} |Q|^4) dx \\ & + D_0 \|(\nabla c, \Delta c)\|_{L^2}^2 + \Gamma \|(\nabla Q, \Delta Q)\|_{L^2}^2 + c_* \Gamma \|Q\|_{L^4}^4 + c_*^2 \Gamma \|Q\|_{L^6}^6 \\ & = (u \cdot \nabla c, \Delta c) - (u \cdot \nabla c, c) + (u \cdot \nabla Q, \Delta Q) - (u \cdot \nabla Q, Q + c_* Q |Q|^2) - (\Omega Q - Q \Omega, \Delta Q) \\ & + (\Omega Q - Q \Omega, Q + c_* Q |Q|^2) + \frac{1}{2} \Gamma ((c - c_*) Q, \Delta Q - Q - c_* Q \operatorname{tr}(Q^2)) \\ & - b \Gamma (Q^2, \Delta Q - Q - c_* Q \operatorname{tr}(Q^2)) + 2c_* \Gamma (Q |Q|^2, \Delta Q) \\ & \leq \frac{D_0}{2} \|\Delta c\|_{L^2}^2 + \frac{\Gamma}{2} \|\Delta Q\|_{L^2}^2 + \frac{c_*^2 \Gamma}{2} \|Q\|_{L^6}^6 + C (\|(c, Q, \nabla c, \nabla Q)\|_{L^2}^2 + \|Q\|_{L^4}^4). \end{aligned}$$

By the Gronwall inequality, we obtain

$$\|(c, Q)\|_{L^\infty(0, T; H^1(\mathcal{O}))} + \|(c, Q)\|_{L^2(0, T; H^2(\mathcal{O}))} \leq C,$$

where $C > 0$ depends on $N, b, c_*, \Gamma, T, \|Q_0\|_{H^1(\mathcal{O})}$, and $\|c_0\|_{H^1(\mathcal{O})}$.

3. For the uniqueness, we denote $(\tilde{c}, \tilde{Q}) = (c_1 - c_2, Q_1 - Q_2)$ for any two solutions (c_1, Q_1) and (c_2, Q_2) of (3.13). Then (\tilde{c}, \tilde{Q}) satisfies

$$\begin{cases} \partial_t \tilde{c} + u \cdot \nabla \tilde{c} = D_0 \Delta \tilde{c}, \\ \partial_t \tilde{Q} + u \cdot \nabla \tilde{Q} - \Omega \tilde{Q} + \tilde{Q} \Omega - \Gamma \Delta \tilde{Q} + \frac{1}{2} \Gamma \tilde{c} Q_1 + \frac{1}{2} \Gamma \tilde{Q} (c_2 - c_*) \\ = \Gamma \left(b(\tilde{Q}(Q_1 + Q_2) - \frac{1}{3} \text{tr}(\tilde{Q}(Q_1 + Q_2)) I_3) - c_* \tilde{Q} \text{tr} Q_1^2 - c_* Q_2 \text{tr}(\tilde{Q}(Q_1 + Q_2)) \right), \end{cases} \quad (3.14)$$

with $(\tilde{c}, \tilde{Q})|_{t=0} = (0, 0)$, $\nabla \tilde{Q} \cdot \vec{n}|_{\partial \mathcal{O}} = 0$, and $\nabla \tilde{c} \cdot \vec{n}|_{\partial \mathcal{O}} = 0$. We sum up the first equation in (3.14) multiplied by \tilde{c} and the second equation in (3.14) multiplied by \tilde{Q} , take the trace, and integrate by parts over \mathcal{O} to obtain

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|(\tilde{c}, \tilde{Q})\|_{L^2}^2 + D_0 \|\nabla \tilde{c}\|_{L^2}^2 + \Gamma \|\nabla \tilde{Q}\|_{L^2}^2 \\ &= -(u \cdot \nabla \tilde{c}, \tilde{c}) - (u \cdot \nabla \tilde{Q}, \tilde{Q}) + (\Omega \tilde{Q} - \tilde{Q} \Omega, \tilde{Q}) - \frac{\Gamma}{2} (\tilde{c} Q_1, \tilde{Q}) - \frac{\Gamma}{2} ((c_2 - c_*) \tilde{Q}, \tilde{Q}) \\ & \quad + \Gamma (b(\tilde{Q}(Q_1 + Q_2) - \frac{1}{3} \text{tr}(\tilde{Q}(Q_1 + Q_2)) I_3) - c_* \tilde{Q} \text{tr} Q_1^2 - c_* Q_2 \text{tr}(\tilde{Q}(Q_1 + Q_2)), \tilde{Q}) \\ &\leq \frac{D_0}{2} \|\nabla \tilde{c}\|_{L^2}^2 + C(\|(\tilde{c}, \tilde{Q})\|_{L^2}^2 + \|\tilde{Q}\|_{L^2} \|\tilde{Q}\|_{L^6} (\|Q_1 + Q_2\|_{L^3} + \|(Q_1, Q_2)\|_{L^6}^2) + \|\nabla \tilde{Q}\|_{L^2} \|\tilde{Q}\|_{L^2}) \\ &\leq \frac{D_0}{2} \|\nabla \tilde{c}\|_{L^2}^2 + \frac{\Gamma}{2} \|\nabla \tilde{Q}\|_{L^2}^2 + C\|(\tilde{c}, \tilde{Q})\|_{L^2}^2, \end{aligned}$$

where C depends on b, c_*, Γ, N , and $\|Q_0\|_{H^1}$, and we have also used Sobolev's embedding inequality, Poincaré inequality, and Young's inequality in the last step. Therefore, applying the Gronwall inequality, we obtain the desired uniqueness result.

4. Now we show that the map: $u \mapsto (c[u], Q[u])$ is continuous. Let $\{u_n\}$ be a bounded sequence in \mathcal{N}_N with

$$\lim_{n \rightarrow \infty} \|u_n - u\|_{C(0, T; C_0^2(\bar{\mathcal{O}}))} = 0,$$

for some $u \in C(0, T; C_0^2(\bar{\mathcal{O}}))$. Denote $(c_n, Q_n) = (c[u_n], Q[u_n])$, $(c, Q) = (c[u], Q[u])$, and $(\tilde{c}_n, \tilde{Q}_n) = (c_n - c, Q_n - Q)$. Taking the difference of the equations satisfied by (c_n, Q_n) and (c, Q) , multiplying the resulting equations by $(-\Delta \tilde{c}_n, -\Delta \tilde{Q}_n)$, taking the trace, and integrating by parts over \mathcal{O} , we have

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|(\nabla \tilde{c}_n, \nabla \tilde{Q}_n)\|_{L^2}^2 + D_0 \|\Delta \tilde{c}_n\|_{L^2}^2 + \Gamma \|\Delta \tilde{Q}_n\|_{L^2}^2 \\ &= (u_n \cdot \nabla \tilde{c}_n + (u_n - u) \cdot \nabla c, \Delta \tilde{c}_n) + (u_n \cdot \nabla \tilde{Q}_n + (u_n - u) \cdot \nabla Q, \Delta \tilde{Q}_n) \\ & \quad - (\Omega_n \tilde{Q}_n + (\Omega_n - \Omega) Q, \Delta \tilde{Q}_n) + (\tilde{Q}_n \Omega_n + Q(\Omega_n - \Omega), \Delta \tilde{Q}_n) \\ & \quad + \frac{\Gamma}{2} (\tilde{c}_n Q_n + (c - c_*) \tilde{Q}_n, \Delta \tilde{Q}_n) - \frac{1}{3} b \Gamma (3 \tilde{Q}_n (Q_n + Q) + \text{tr}(\tilde{Q}_n (Q_n + Q)) I_3, \Delta \tilde{Q}_n) \\ & \quad + c_* \Gamma (|Q_n|^2 \tilde{Q}_n + Q(|Q_n|^2 - |Q|^2), \Delta \tilde{Q}_n) \\ &= \sum_{i=1}^7 \mathcal{I}_i. \end{aligned}$$

In the following, we estimate all the terms on the right-hand side of the above equation:

$$\begin{aligned} |\mathcal{I}_1| &= |(u_n \cdot \nabla \tilde{c}_n + (u_n - u) \cdot \nabla c, \Delta \tilde{c}_n)| \\ &\leq C(\|u_n\|_{L^\infty} \|\nabla \tilde{c}_n\|_{L^2} + \|u_n - u\|_{L^\infty} \|\nabla c\|_{L^2}) \|\Delta \tilde{c}_n\|_{L^2} \\ &\leq \frac{D_0}{2} \|\Delta \tilde{c}_n\|_{L^2}^2 + C \|\nabla \tilde{c}_n\|_{L^2}^2 + C \|u_n - u\|_{L^\infty}^2, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_2| &= |(u_n \cdot \nabla \tilde{Q}_n + (u_n - u) \cdot \nabla Q, \Delta \tilde{Q}_n)| \\ &\leq C(\|u_n\|_{L^\infty} \|\nabla \tilde{Q}_n\|_{L^2} + \|u_n - u\|_{L^\infty} \|\nabla Q\|_{L^2}) \|\Delta \tilde{Q}_n\|_{L^2} \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\nabla \tilde{Q}_n\|_{L^2}^2 + C \|u_n - u\|_{L^\infty}^2, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_3| &= |(\Omega_n \tilde{Q}_n + (\Omega_n - \Omega)Q, \Delta \tilde{Q}_n)| \\ &\leq C \|\Omega_n\|_{L^\infty} \|\tilde{Q}_n\|_{L^2} \|\Delta \tilde{Q}_n\|_{L^2} + \|\nabla u_n - \nabla u\|_{L^\infty} \|Q\|_{L^2} \|\Delta \tilde{Q}_n\|_{L^2} \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\tilde{Q}_n\|_{L^2}^2 + C \|\nabla u_n - \nabla u\|_{L^\infty}^2 \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\nabla \tilde{Q}_n\|_{L^2}^2 + C \|\nabla u_n - \nabla u\|_{L^\infty}^2, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_4| &= |(\tilde{Q}_n \Omega_n + Q(\Omega_n - \Omega), \Delta \tilde{Q}_n)| \\ &\leq C \|\tilde{Q}_n\|_{L^2} \|\Omega_n\|_{L^\infty} \|\Delta \tilde{Q}_n\|_{L^2} + \|Q\|_{L^2} \|\Omega_n - \Omega\|_{L^\infty} \|\Delta \tilde{Q}_n\|_{L^2} \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\tilde{Q}_n\|_{L^2}^2 + C \|\nabla u_n - \nabla u\|_{L^\infty}^2 \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\nabla \tilde{Q}_n\|_{L^2}^2 + C \|\nabla u_n - \nabla u\|_{L^\infty}^2, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_5| &= \frac{\Gamma}{2} |(\tilde{c}_n Q_n + (c - c_*) \tilde{Q}_n, \Delta \tilde{Q}_n)| \leq C(\|\tilde{c}_n\|_{L^4} \|Q_n\|_{L^4} + \|c - c_*\|_{L^\infty} \|\tilde{Q}_n\|_{L^2}) \|\Delta \tilde{Q}_n\|_{L^2} \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\nabla \tilde{c}_n\|_{L^2}^2 + C \|\tilde{Q}_n\|_{L^2}^2, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_6| &= \frac{1}{3} b \Gamma |(3 \tilde{Q}_n(Q_n + Q) + \text{tr}(\tilde{Q}_n(Q_n + Q)) \mathbf{I}_3, \Delta \tilde{Q}_n)| \leq C \|\tilde{Q}_n\|_{L^3} \|Q_n + Q\|_{L^6} \|\Delta \tilde{Q}_n\|_{L^2} \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\nabla \tilde{Q}_n\|_{L^2}^2, \end{aligned}$$

and

$$\begin{aligned} |\mathcal{I}_7| &= c_* \Gamma (|Q_n|^2 \tilde{Q}_n + Q(|Q_n|^2 - |Q|^2), \Delta \tilde{Q}_n) \\ &\leq C \Gamma (\|Q_n\|_{L^6}^2 + \|Q_n\|_{L^6} \|Q\|_{L^6} + \|Q\|_{L^6}^2) \|\tilde{Q}_n\|_{L^6} \|\Delta \tilde{Q}_n\|_{L^2} \\ &\leq \frac{\Gamma}{12} \|\Delta \tilde{Q}_n\|_{L^2}^2 + C \|\nabla \tilde{Q}_n\|_{L^2}^2. \end{aligned}$$

Combining all the above estimates, we conclude

$$\frac{1}{2} \frac{d}{dt} \|(\nabla \tilde{c}_n, \nabla \tilde{Q}_n)\|_{L^2}^2 + \frac{D_0}{2} \|\Delta \tilde{c}_n\|_{L^2}^2 + \frac{\Gamma}{2} \|\Delta \tilde{Q}_n\|_{L^2}^2$$

$$\leq C\|u_n - u\|_{C_0^2(\mathcal{O})}^2 + C\|(\nabla \tilde{c}_n, \nabla \tilde{Q}_n)\|_{L^2}^2.$$

Then, by the Gronwall inequality, we have

$$\begin{aligned} & \|(\nabla \tilde{c}_n(t, \cdot), \nabla \tilde{Q}_n(t, \cdot))\|_{L^2}^2 + \Gamma \int_0^t \|(\Delta \tilde{c}_n(s, \cdot), \Delta \tilde{Q}_n(s, \cdot))\|_{L^2}^2 ds \\ & \leq CT e^{CT} \|u_n - u\|_{L^\infty(0, T; C_0^2(\mathcal{O}))}^2. \end{aligned} \quad (3.15)$$

As $n \rightarrow \infty$, we conclude

$$\lim_{n \rightarrow \infty} (\|(\tilde{c}_n, \tilde{Q}_n)\|_{L^\infty(0, T; H^1)} + \|(\tilde{c}_n, \tilde{Q}_n)\|_{L^2(0, T; H^2)}) = 0.$$

5. Finally, we show that $Q \in S_0^3$, i.e., $Q^\top = Q$ and $\text{tr } Q = 0$ a.e. in \mathcal{O}_T . It is clear that, if Q is a solution of problem (3.13), so is Q^\top . Then, by the uniqueness of the solution, we know that $Q^\top = Q$.

Then the only thing left is to show that $\text{tr } Q = 0$. We take the trace on both sides of the second equation in (3.13) to obtain

$$\partial_t(\text{tr } Q) + u \cdot \nabla \text{tr } Q = \Gamma(\Delta \text{tr } Q - \frac{1}{2}(c - c_*)\text{tr } Q - c_* \text{tr } Q \text{tr } Q^2), \quad (3.16)$$

with $\text{tr } Q|_{t=0} = \text{tr } Q_0 = 0$ and $\nabla \text{tr } Q \cdot \vec{n}|_{\partial \mathcal{O}} = 0$, where we have used the fact that $Q^\top = Q$, $\Omega^\top = -\Omega$, and $\text{tr}(Q\Omega) = \text{tr}(\Omega Q)$. Then we multiply equation (3.16) by $\text{tr } Q$ and integrate by parts over \mathcal{O} to obtain

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\text{tr } Q\|_{L^2}^2 + \Gamma \|\nabla \text{tr } Q\|_{L^2}^2 \\ & = -\frac{1}{2} \Gamma ((c - c_*)\text{tr } Q, \text{tr } Q) - c_* \Gamma \int_{\mathcal{O}} |\text{tr } Q|^2 \text{tr } Q^2 dx - \int_{\mathcal{O}} u \cdot \nabla (\text{tr } Q) \text{tr } Q dx \\ & \leq C \|\text{tr } Q\|_{L^2}^2 + C \|Q\|_{L^6}^2 \|\text{tr } Q\|_{L^6} \|\text{tr } Q\|_{L^2} + C \|\nabla \text{tr } Q\|_{L^2} \|\text{tr } Q\|_{L^2} \\ & \leq C \|\text{tr } Q\|_{L^2}^2 + C \|\nabla \text{tr } Q\|_{L^2} \|\text{tr } Q\|_{L^2} + C \|\nabla \text{tr } Q\|_{L^2} \|\text{tr } Q\|_{L^2} \\ & \leq \frac{\Gamma}{2} \|\nabla \text{tr } Q\|_{L^2}^2 + C \|\text{tr } Q\|_{L^2}^2. \end{aligned}$$

Applying the Gronwall inequality again, we complete the proof. \square

3.2. The Faedo-Galerkin approximation scheme. In this section, we proceed to solve (3.3) by the Faedo-Galerkin approximation scheme. Let $\{\psi_n\}$ be a family of smooth eigenfunctions of the Laplacian operator:

$$\begin{aligned} -\Delta \psi_n &= \lambda_n \psi_n & \text{on } \mathcal{O}, \\ \psi_n|_{\partial \mathcal{O}} &= 0, \end{aligned}$$

where $0 < \lambda_1 \leq \lambda_2 \leq \dots$ are eigenvalues. We know that the eigenfunctions $\{\psi_n\}_{n=1}^\infty$ form an orthogonal basis of $H_0^1(\mathcal{O})$.

Now, consider a sequence of finite-dimensional spaces:

$$X_n = \text{span}\{\psi_1, \psi_2, \dots, \psi_n\}, \quad n = 1, 2, \dots, \quad (3.17)$$

and look for solutions $u_n \in C(0, T; X_n)$ to the following variational approximation problem:

$$\begin{aligned} & \int_{\mathcal{O}} \rho(t, x) u_n(t, x) \cdot \psi(x) \, dx - \int_{\mathcal{O}} m_0(x) \cdot \psi(x) \, dx \\ &= \int_0^t \int_{\mathcal{O}} (\mu \Delta u_n + (\mu + \nu) \nabla \operatorname{div} u_n - \operatorname{div}(\rho u_n \otimes u_n) - \nabla(\rho^\gamma + \delta \rho^\beta) - \varepsilon \nabla \rho \cdot \nabla u_n) \cdot \psi \, dx \, ds \\ &+ \int_0^t \int_{\mathcal{O}} \nabla \cdot (F(Q) \mathbf{I}_3 - \nabla Q \odot \nabla Q + Q \Delta Q - \Delta Q Q + \sigma_* c^2 Q) \cdot \psi \, dx \, ds \end{aligned} \quad (3.18)$$

for any $t \in [0, T]$, and $\psi \in X_n$.

Next we introduce a family of operators, as in [15]:

$$\mathcal{M}[\rho] : X_n \mapsto X_n^*, \quad (\mathcal{M}[\rho]v, w) = \int_{\mathcal{O}} \rho v \cdot w \, dx \quad \text{for any } v, w \in X_n.$$

The map:

$$\rho \mapsto \mathcal{M}^{-1}[\rho]$$

from $N_\eta = \{\rho \in L^1(\mathcal{O}) : \inf_{x \in \Omega} \rho \geq \eta > 0\}$ into $\mathcal{L}(X_n^*, X_n)$ has the following property:

$$\|\mathcal{M}^{-1}[\rho_1] - \mathcal{M}^{-1}[\rho_2]\|_{\mathcal{L}(X_n^*, X_n)} \leq C(n, \eta) \|\rho_1 - \rho_2\|_{L^1(\mathcal{O})}. \quad (3.19)$$

Using Theorems 3.1–3.2 with $\rho_n = S[u_n]$, $c_n = c[u_n]$, and $Q_n = Q[u_n]$, we can rewrite the variational problem (3.18) as

$$u_n(t) = \mathcal{M}^{-1}[\rho_n] \left(m^* + \int_0^t \mathcal{N}[c_n(s), \rho_n(s), u_n(s), Q_n(s)] \, ds \right)$$

with

$$\begin{aligned} (m^*, \psi) &= \int_{\mathcal{O}} m_0 \cdot \psi \, dx, \\ (\mathcal{N}[c_n, \rho_n, u_n, Q_n], \psi) &= \int_{\mathcal{O}} (\mu \Delta u_n + (\mu + \nu) \nabla \operatorname{div} u_n - \operatorname{div}(\rho_n u_n \otimes u_n) - \nabla(\rho_n^\gamma + \delta \rho_n^\beta)) \cdot \psi \, dx \\ &- \varepsilon \int_{\mathcal{O}} (\nabla \rho_n \cdot \nabla) u_n \cdot \psi \, dx + \int_{\mathcal{O}} \nabla \cdot (-\nabla Q_n \odot \nabla Q_n + F(Q_n) \mathbf{I}_3) \cdot \psi \, dx \\ &+ \int_{\mathcal{O}} \nabla \cdot (Q_n \Delta Q_n - \Delta Q_n Q_n + \sigma_* c_n^2 Q_n) \cdot \psi \, dx, \end{aligned}$$

for any $t \in [0, T]$ and $\psi \in X_n$. Therefore, combining (3.12), (3.15), and (3.19), we achieve a local solution (c_n, ρ_n, u_n, Q_n) of problem (3.1)–(3.2), (3.4), and (3.18), with initial-boundary data (3.6)–(3.10) on a short time interval $[0, T_n]$, $T_n \leq T$, by using the standard fixed point theorem on $C(0, T; X_n)$. In order to extend the existence time T_n to T for any $n = 1, 2, \dots$, we need to prove that u_n stays bounded in X_n for the whole interval $[0, T_n]$. Hence, in the following, we establish an energy inequality as in Proposition 2.1.

Differentiate (3.18) with respect to t and take $\psi = u_n$ as a test function to obtain

$$\begin{aligned} & \frac{d}{dt} \int_{\mathcal{O}} \left(\frac{1}{2} \rho_n |u_n|^2 + \frac{\rho_n^\gamma}{\gamma-1} + \frac{\delta \rho_n^\beta}{\beta-1} \right) dx + \mu \|\nabla u_n\|_{L^2}^2 + (\nu + \mu) \|\operatorname{div} u_n\|_{L^2}^2 \\ & \quad + \varepsilon \int_{\mathcal{O}} (\gamma \rho_n^{\gamma-2} + \delta \beta \rho_n^{\beta-2}) |\nabla \rho_n|^2 dx \\ & = \int_{\mathcal{O}} \nabla \cdot (-\nabla Q_n \odot \nabla Q_n + F(Q_n) I_3 + Q_n \Delta Q_n - \Delta Q_n Q_n + \sigma_* c_n^2 Q_n) \cdot u_n dx, \end{aligned} \quad (3.20)$$

where we have used the following equalities as in Proposition 2.1:

$$\begin{aligned} & \int_{\mathcal{O}} \partial_t (\rho_n u_n) \cdot u_n dx + \int_{\mathcal{O}} \operatorname{div} (\rho_n u_n \otimes u_n) \cdot u_n dx \\ & = \frac{1}{2} \frac{d}{dt} \int_{\mathcal{O}} \rho_n |u_n|^2 dx - \varepsilon \int_{\mathcal{O}} (\nabla \rho_n \cdot \nabla) u_n \cdot u_n dx, \\ & \int_{\mathcal{O}} \nabla \rho_n^\gamma \cdot u_n dx = \frac{d}{dt} \int_{\mathcal{O}} \frac{\rho_n^\gamma}{\gamma-1} dx + \varepsilon \gamma \int_{\mathcal{O}} \rho_n^{\gamma-2} |\nabla \rho_n|^2 dx, \\ & \int_{\mathcal{O}} \delta \nabla \rho_n^\beta \cdot u_n dx = \frac{d}{dt} \int_{\mathcal{O}} \frac{\delta \rho_n^\beta}{\beta-1} dx + \varepsilon \delta \beta \int_{\mathcal{O}} \rho_n^{\beta-2} |\nabla \rho_n|^2 dx. \end{aligned}$$

Then we take the inner product of (3.1) with c_n , (3.4) with $-(\Delta Q_n - Q_n - c_* Q_n \operatorname{tr}(Q_n^2))$, add the resulting equations to (3.20) to obtain

$$\begin{aligned} & \frac{d}{dt} E_\delta^n(t) + D_0 \|\nabla c_n\|_{L^2}^2 + \mu \|\nabla u_n\|_{L^2}^2 + (\mu + \nu) \|\operatorname{div} u_n\|_{L^2}^2 + \Gamma \|(\nabla Q_n, \Delta Q_n)\|_{L^2}^2 \\ & \quad + c_* \Gamma \|Q_n\|_{L^4}^4 + c_*^2 \Gamma \|Q_n\|_{L^6}^6 + \varepsilon \int_{\mathcal{O}} (\gamma \rho_n^{\gamma-2} + \delta \beta \rho_n^{\beta-2}) |\nabla \rho_n|^2 dx \\ & = \int_{\mathcal{O}} \nabla \cdot (-\nabla Q_n \odot \nabla Q_n + F(Q_n) I_3 + Q_n \Delta Q_n - \Delta Q_n Q_n + \sigma_* c_n^2 Q_n) \cdot u_n dx \\ & \quad + (u_n \cdot \nabla Q_n, \Delta Q_n) - (u_n \cdot \nabla Q_n, Q_n + c_* Q_n |Q_n|^2) - (\Omega_n Q_n - Q_n \Omega_n, \Delta Q_n) \\ & \quad + (\Omega_n Q_n - Q_n \Omega_n, Q_n + c_* Q_n |Q_n|^2) + \frac{1}{2} \Gamma ((c_n - c_*) Q_n, \Delta Q_n - Q_n - c_* Q_n \operatorname{tr}(Q_n^2)) \\ & \quad - b \Gamma (Q_n^2, \Delta Q_n - Q_n - c_* Q_n \operatorname{tr}(Q_n^2)) + 2 c_* \Gamma (Q_n |Q_n|^2, \Delta Q_n) - (u_n \cdot \nabla c_n, c_n) \\ & \leq \frac{D_0}{2} \|\nabla c_n\|_{L^2}^2 + \frac{\mu}{2} \|\nabla u_n\|_{L^2}^2 + \frac{\Gamma}{2} \|\Delta Q_n\|_{L^2}^2 + \frac{c_* \Gamma}{2} \|Q_n\|_{L^6}^6 + C(\|(u_n, Q_n)\|_{L^2}^2 + \|Q_n\|_{L^4}^4) \end{aligned}$$

where C is independent of n and ε , and

$$E_\delta^n(t) = \int_{\mathcal{O}} \left(\frac{1}{2} |c_n|^2 + \frac{1}{2} \rho_n |u_n|^2 + \frac{\rho_n^\gamma}{\gamma-1} + \frac{\delta \rho_n^\beta}{\beta-1} + F(Q_n) \right) dx.$$

This implies

$$\begin{aligned} & \frac{d}{dt} E_\delta^n(t) + \frac{D_0}{2} \|\nabla c_n\|_{L^2}^2 + \frac{\mu}{2} \|\nabla u_n\|_{L^2}^2 + (\mu + \nu) \|\operatorname{div} u_n\|_{L^2}^2 + \Gamma \|\nabla Q_n\|_{L^2}^2 \\ & \quad + \frac{\Gamma}{2} \|\Delta Q_n\|_{L^2}^2 + \frac{c_*^2 \Gamma}{2} \|Q_n\|_{L^6}^6 + \varepsilon \int_{\mathcal{O}} (\gamma \rho_n^{\gamma-2} + \delta \beta \rho_n^{\beta-2}) |\nabla \rho_n|^2 dx \\ & \leq C(\|(u_n, Q_n)\|_{L^2}^2 + \|Q_n\|_{L^4}^4). \end{aligned} \quad (3.21)$$

The above inequality yields

$$\mu \int_0^{T_n} \|\nabla u_n\|_{L^2}^2 dt \leq C, \quad (3.22)$$

$$\sup_{t \in [0, T_n]} \int_{\mathcal{O}} \rho_n |u_n|^2 dx \leq C, \quad (3.23)$$

where C is a constant independent of n . Since X_n is a finite-dimension space, we can deduce from Lemma 3.1 that there exists a constant $C = C(n, c_0, \rho_0, m_0, Q_0, b, \mathcal{O})$ such that

$$0 < C \leq \rho_n(t, x) \leq \frac{1}{C} \quad \text{for all } t \in (0, T_n) \text{ and } x \in \mathcal{O}, \quad (3.24)$$

which, combined with (3.23) and the fact that the L^∞ and L^2 norms are equivalent on X_n , yields

$$\sup_{t \in [0, T_n]} \|(u_n, \nabla u_n)(t, \cdot)\|_{L^\infty(\mathcal{O})} \leq C(n, E_\delta^n(0), N, \mathcal{O}).$$

Then we can extend the existence time-interval $[0, T_n]$ of (c_n, ρ_n, u_n, Q_n) to $[0, T]$.

We summarize the results in this subsection in the following lemma.

Lemma 3.3. *Let $\beta \geq 4$. Then there exists a solution (c_n, ρ_n, u_n, Q_n) of problem (3.1)–(3.2), (3.4), and (3.18) with the corresponding initial-boundary data (3.5)–(3.10). Moreover, the following estimates hold:*

$$0 < \underline{c} \leq c_n(t, x) \leq \bar{c}, \quad \|c_n\|_{L^\infty(0, T; L^2(\mathcal{O})) \cap L^2(0, T; H^1(\mathcal{O}))} \leq C, \quad (3.25)$$

$$\sup_{t \in [0, T]} \|\rho_n(t, \cdot)\|_{L^\gamma(\mathcal{O})}^\gamma \leq C, \quad (3.26)$$

$$\delta \sup_{t \in [0, T]} \|\rho_n(t, \cdot)\|_{L^\beta(\mathcal{O})}^\beta \leq C, \quad (3.27)$$

$$\varepsilon \|\nabla \rho_n^{\frac{\gamma}{2}}\|_{L^2(0, T; L^2(\mathcal{O}))}^2 + \varepsilon \delta \|\nabla \rho_n^{\frac{\beta}{2}}\|_{L^2(0, T; L^2(\mathcal{O}))}^2 \leq C, \quad (3.28)$$

$$\sup_{t \in [0, T]} \|\sqrt{\rho_n}(t, \cdot) u_n(t, \cdot)\|_{L^2(\mathcal{O})}^2 \leq C, \quad (3.29)$$

$$\|u_n\|_{L^2(0, T; H_0^1(\mathcal{O}))} \leq C, \quad (3.30)$$

$$\|\rho_n\|_{L^{\beta+1}(\mathcal{O}_T)} \leq C, \quad (3.31)$$

$$\varepsilon \|\nabla \rho_n\|_{L^2(0, T; L^2(\mathcal{O}))}^2 \leq C, \quad (3.32)$$

$$\|Q_n\|_{L^\infty(0, T; H^1(\mathcal{O}) \cap L^4(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}) \cap L^6(\mathcal{O}))} \leq C, \quad (3.33)$$

$$\|Q_n\|_{L^{10}(\mathcal{O}_T)} \leq C, \quad (3.34)$$

$$\|\nabla Q_n\|_{L^{\frac{10}{3}}(\mathcal{O}_T)} \leq C, \quad (3.35)$$

where C is a constant independent of n and ε .

Proof. Estimates (3.25)–(3.30) and (3.33) follow from the energy estimate (3.21). Moreover, we can use similar methods to obtain (3.34)–(3.35) as in Corollary 2.1. We only need to show (3.31)–(3.32).

From (3.28), $\rho_n^{\frac{\beta}{2}} \in L_t^2 H_x^1$. This, together with the embedding: $H^1(\mathcal{O}) \subset L^6(\mathcal{O})$, yields

$$\|\rho_n^\beta\|_{L_t^1 L_x^3} \leq C \quad (3.36)$$

with C independent of n . Combining (3.27), (3.36), and the interpolation (pp. 623, [13]):

$$\|\rho_n^\beta\|_{L_x^2} \leq C \|\rho_n^\beta\|_{L_x^1}^{\frac{1}{4}} \|\rho_n^\beta\|_{L_x^3}^{\frac{3}{4}},$$

we have

$$\|\rho_n^\beta\|_{L_t^{\frac{4}{3}} L_x^2}^{\frac{4}{3}} \leq C \int_0^T \|\rho_n^\beta\|_{L_x^1}^{\frac{1}{3}} \|\rho_n^\beta\|_{L_x^3} dt \leq C \|\rho_n^\beta\|_{L_t^\infty L_x^1}^{\frac{1}{3}} \|\rho_n^\beta\|_{L_t^1 L_x^3} \leq C,$$

which implies that $\rho_n \in L^{\frac{4}{3}\beta}(0, T; L^{2\beta}(\mathcal{O}))$. Moreover, this, together with the following interpolation:

$$\|\rho_n\|_{L_x^{\frac{4}{3}\beta}} \leq C \|\rho_n\|_{L_x^\beta}^{\frac{1}{2}} \|\rho_n\|_{L_x^{2\beta}}^{\frac{1}{2}},$$

gives

$$\|\rho_n\|_{L_{t,x}^{\frac{4\beta}{3}}}^{\frac{4\beta}{3}} \leq C \int_0^T \|\rho_n\|_{L_x^\beta}^{\frac{2}{3}\beta} \|\rho_n\|_{L_x^{2\beta}}^{\frac{2}{3}\beta} dt \leq C(\delta),$$

i.e., $\rho_n \in L^{\frac{4}{3}\beta}(\mathcal{O}_T)$. Then, if $\beta \geq 3$ ($\Leftrightarrow \frac{4}{3}\beta \geq \beta + 1$), we have

$$\|\rho_n\|_{L^{\beta+1}(\mathcal{O}_T)} \leq C(\delta).$$

Regarding (3.32), we multiply (3.2) by ρ_n and integrate by parts over \mathcal{O} to obtain

$$\frac{1}{2} \frac{d}{dt} \|\rho_n\|_{L_x^2}^2 + \varepsilon \|\nabla \rho_n\|_{L_x^2}^2 = -\frac{1}{2} \int_{\mathcal{O}} \operatorname{div} u_n \rho_n^2 dx \leq C \|\operatorname{div} u_n\|_{L_x^2} \|\rho_n\|_{L_x^4}^2,$$

which yields

$$\varepsilon \|\nabla \rho_n\|_{L^2(\mathcal{O}_T)}^2 \leq \frac{1}{2} (\|\rho_0\|_{L_x^2}^2 + \sqrt{T} \|\rho_n\|_{L_t^\infty L_x^4}^2 \|\operatorname{div} u_n\|_{L_{t,x}^2}).$$

Therefore, (3.32) follows from (3.27) and (3.30), provided $\beta \geq 4$. \square

3.3. The existence of the first level approximate solutions. In this subsection, we obtain a solution (c, ρ, u, Q) of problem (3.1)–(3.10), by letting $n \rightarrow \infty$. We do not distinguish between the sequence convergence and the subsequence convergence for the sake of convenience. Assume that $\beta > 4$ and $\gamma > \frac{3}{2}$. It follows from [15] that, as $n \rightarrow \infty$,

$$\rho_n \rightarrow \rho \quad \text{in } L^4(\mathcal{O}_T), \quad (3.37)$$

$$\rho_n^\gamma \rightarrow \rho^\gamma, \quad \rho_n^\beta \rightarrow \rho^\beta \quad \text{in } L^1(\mathcal{O}_T) \text{ if } \beta > \gamma, \quad (3.38)$$

$$u_n \rightharpoonup u \quad \text{in } L^2(0, T; H_0^1(\mathcal{O})). \quad (3.39)$$

Moreover, we can infer that $\partial_t c_n \in L^1(0, T; H^{-1}(\mathcal{O}))$ from (3.25), (3.30), and the fact that $\partial_t c_n$ satisfies equation (3.1). Applying the Aubin-Lions lemma, we have

$$c_n \rightharpoonup c \quad \text{in } L^2(0, T; H^1(\mathcal{O})), \quad c_n \rightarrow c \quad \text{in } L^2(0, T; L^2(\mathcal{O})). \quad (3.40)$$

Thus, we know that the limit function c is a weak solution to (3.1). Similarly, estimates (3.25), (3.30), and (3.33), along with the fact that $\partial_t Q_n$ satisfies (3.4), yield that $\partial_t Q_n \in L^2(0, T; L^{\frac{3}{2}}(\mathcal{O}))$. In fact,

$$\begin{aligned} \|(u_n \cdot \nabla) Q_n\|_{L_t^2 L_x^{\frac{3}{2}}} &\leq C \left(\int_0^T \|u_n\|_{L_x^6}^2 \|\nabla Q_n\|_{L_x^2}^2 dt \right)^{\frac{1}{2}} \leq C \|u_n\|_{L_t^2(H_0^1)_x} \|\nabla Q_n\|_{L_t^\infty L_x^2} \leq C, \\ \|Q_n \Omega_n - \Omega_n Q_n\|_{L_t^2 L_x^{\frac{3}{2}}} &\leq C \left(\int_0^T \|Q_n\|_{L_x^6}^2 \|\nabla u_n\|_{L_x^2}^2 dt \right)^{\frac{1}{2}} \leq C \|Q_n\|_{L_t^\infty H_x^1} \|\nabla u_n\|_{L^2(\mathcal{O}_T)} \leq C, \\ \|\Gamma H[Q_n, c_n]\|_{L_t^2 L_x^{\frac{3}{2}}} &= \Gamma \left\| \Delta Q_n - \frac{c_n - c_*}{2} Q_n + b[Q_n^2 - \frac{\text{tr}(Q_n^2)}{3} \text{Id}] - c_* Q_n |Q_n|^2 \right\|_{L_t^2 L_x^{\frac{3}{2}}} \\ &\leq C \left(\|(Q_n, \Delta Q_n)\|_{L^2(\mathcal{O}_T)} + (\|c_n\|_{L_t^2 H_x^1} + \|Q_n\|_{L_t^2 H_x^2}) (\|Q_n\|_{L_t^\infty H_x^1} + \|Q_n\|_{L_t^\infty H_x^1}^2) \right) \leq C. \end{aligned}$$

Then, by the Aubin-Lions lemma, we have

$$Q_n \rightharpoonup Q \text{ in } L^2(0, T; H^2(\mathcal{O})), \quad Q_n \rightarrow Q \text{ in } L^2(0, T; H^1(\mathcal{O})). \quad (3.41)$$

This ensures that we can pass to the limit in equation (3.4) in $\mathcal{D}'(\mathcal{O}_T)$ as $n \rightarrow \infty$, i.e., Q is a weak solution of (3.4).

Furthermore, we also see that $\rho_n u_n$ is bounded in $L^\infty(0, T; L^{\frac{2\gamma}{\gamma+1}}(\mathcal{O}))$ with $\frac{2\gamma}{\gamma+1} > \frac{6}{5}$ (since $\gamma > \frac{3}{2}$), by using (3.26) and (3.29)–(3.30). In fact,

$$\|\rho_n u_n\|_{L_x^{\frac{2\gamma}{\gamma+1}}} \leq C \|\rho_n\|_{L_x^{\frac{1}{\gamma}}}^{\frac{1}{2}} \|\sqrt{\rho_n} u_n\|_{L_x^2}. \quad (3.42)$$

This, together with (3.37) and (3.39), yields

$$\rho_n u_n \xrightarrow{*} \rho u \text{ in } L^\infty(0, T; L^{\frac{2\gamma}{\gamma+1}}(\mathcal{O})).$$

Then we can conclude that ρ is a weak solution of (3.2) and we can pass to the limit in (3.2) as $n \rightarrow \infty$.

In the following, we show that the limit function u satisfies equation (3.18), by using Corollary A.1 in Appendix A. Then we need to establish convergence results for the terms: $\rho_n u_n \otimes u_n$ and $\nabla \rho_n \cdot \nabla u_n$, which require more estimates for the density. From Lemma 2.4 in [15], we know that (3.2) holds in the following strong sense:

Lemma 3.4. *There exist $r > 1$ and $q > 2$ such that*

$$\begin{aligned} \partial_t \rho_n, \Delta \rho_n &\text{ are bounded in } L^r(\mathcal{O}_T), \\ \nabla \rho_n &\text{ is bounded in } L^q(0, T; L^2(\mathcal{O})), \end{aligned}$$

uniformly with respect to n . Consequently, the limit function ρ belongs to the same class and satisfies equation (3.2) almost everywhere on \mathcal{O}_T and the boundary conditions (3.9) in the sense of traces.

To continue the proof, we first show that $\int_{\mathcal{O}} \rho_n u_n \cdot \psi dx$ is equi-continuous in t , for any fixed test function $\psi \in X_n$ in (3.18). Using Lemmas 3.3–3.4, we see that, for any $0 < \xi < 1$,

$$\int_t^{t+\xi} \int_{\mathcal{O}} (\mu \Delta u_n + (\mu + \nu) \nabla \text{div} u_n) \cdot \psi dx ds \leq C \int_t^{t+\xi} \|\nabla u_n\|_{L_x^2} \|\nabla \psi\|_{L_x^2} ds \leq C \sqrt{\xi},$$

$$\begin{aligned}
& \int_t^{t+\xi} \int_{\mathcal{O}} \operatorname{div}(\rho_n u_n \otimes u_n) \cdot \psi \, dx ds \leq \int_t^{t+\xi} \|\sqrt{\rho_n} u_n\|_{L_x^2}^2 \|\nabla \psi\|_{L_x^\infty} \, ds \leq C\xi, \\
& \int_t^{t+\xi} \int_{\mathcal{O}} \nabla(\rho_n^\gamma + \delta \rho_n^\beta) \cdot \psi \, dx ds \leq \int_t^{t+\xi} (\|\rho_n\|_{L_x^\gamma}^\gamma + \delta \|\rho_n\|_{L_x^\beta}^\beta) \|\operatorname{div} \psi\|_{L_x^\infty} \, ds \leq C\xi, \\
& \varepsilon \int_t^{t+\xi} \int_{\mathcal{O}} \nabla \rho_n \cdot \nabla u_n \cdot \psi \, dx ds \leq C\varepsilon \|\nabla \rho_n\|_{L_t^q L_x^2} \|\nabla u_n\|_{L_t^2 L_x^2} \xi^{\frac{1}{2}-\frac{1}{q}} \leq C\varepsilon \xi^{\frac{1}{2}-\frac{1}{q}} \quad \text{for } q > 2, \\
& \int_t^{t+\xi} \int_{\mathcal{O}} \nabla \cdot F(Q_n) \mathbf{I}_3 \cdot \psi \, dx ds = -\frac{1}{2} \int_0^{t+\xi} \int_{\mathcal{O}} (|(Q_n, \nabla Q_n)|^2 + \frac{c_*}{2} |Q_n|^4) \operatorname{div} \psi \, dx ds \\
& \leq C \int_t^{t+\xi} (\|(Q_n, \nabla Q_n)\|_{L_x^2}^2 + \|Q_n\|_{L_x^4}^4) \|\operatorname{div} \psi\|_{L_x^\infty} \, ds \leq C\xi, \\
& \int_t^{t+\xi} \int_{\mathcal{O}} \nabla \cdot (-\nabla Q_n \odot \nabla Q_n + Q_n \Delta Q_n - \Delta Q_n Q_n + \sigma_* c_n^2 Q_n) \cdot \psi \, dx ds \\
& \leq C \int_t^{t+\xi} ((\|\nabla Q_n\|_{L_x^2}^2 + \|Q_n\|_{L_x^2} \|\Delta Q_n\|_{L_x^2}) \|\nabla \psi\|_{L_x^\infty} + \|c_n\|_{L^\infty}^2 \|Q_n\|_{L_x^2} \|\nabla \psi\|_{L_x^2}) \, ds \\
& \leq C(\xi + \sqrt{\xi}).
\end{aligned}$$

Together with the fact that $\rho_n u_n$ is uniformly bounded in $L^\infty(0, T; L^{\frac{2\gamma}{\gamma+1}}(\mathcal{O}))$ with respect to n , and X_n is dense in $L^{\frac{\gamma-1}{2\gamma}}(\mathcal{O})$, we conclude by Corollary A.1 that

$$\rho_n u_n \rightarrow \rho u \quad \text{in } C([0, T]; L_{\text{weak}}^{\frac{2\gamma}{\gamma+1}}(\mathcal{O})) \quad \text{as } n \rightarrow \infty. \quad (3.43)$$

Since $\frac{2\gamma}{\gamma+1} > \frac{6}{5}$ ($\gamma > \frac{3}{2}$), by Proposition A.1, (3.43) yields

$$\rho_n u_n \rightarrow \rho u \quad \text{in } C([0, T]; H^{-1}(\mathcal{O})).$$

This, combined with (3.39), yields

$$\rho_n u_n \otimes u_n \rightarrow \rho u \otimes u \quad \text{in } \mathcal{D}'(\mathcal{O}_T). \quad (3.44)$$

Next, let us elaborate on the convergence result for the term: $\nabla u_n \cdot \nabla \rho_n$. We multiply (3.2) by ρ_n and integrate by parts to obtain

$$\|\rho_n(t, \cdot)\|_{L_x^2}^2 + 2\varepsilon \int_0^t \|\nabla \rho_n(t, \cdot)\|_{L_x^2}^2 \, ds = - \int_0^t \int_{\mathcal{O}} \operatorname{div} u_n |\rho_n|^2 \, dx ds + \|\rho_0\|_{L_x^2}^2. \quad (3.45)$$

By Lemma 3.4, we know that the limit function ρ also satisfies (3.2). Applying the same argument to ρ as above, we have

$$\|\rho(t, \cdot)\|_{L_x^2}^2 + 2\varepsilon \int_0^t \|\nabla \rho(t, \cdot)\|_{L_x^2}^2 \, ds = - \int_0^t \int_{\mathcal{O}} \operatorname{div} u |\rho|^2 \, dx ds + \|\rho_0\|_{L_x^2}^2. \quad (3.46)$$

Differentiating (3.45) with respect to t , we use (3.27), (3.30), and Lemma 3.4 to obtain

$$\frac{d}{dt} \|\rho_n(t, \cdot)\|_{L_x^2}^2 = -2\varepsilon \|\nabla \rho_n(t, \cdot)\|_{L_x^2}^2 - \int_{\mathcal{O}} \operatorname{div} u_n |\rho_n|^2 \, dx + \|\rho_0\|_{L_x^2}^2 \in L^q(0, T), \quad 1 < q < 2,$$

which implies that $\|\rho_n(t, \cdot)\|_{L_x^2}^2$ is equi-continuous. Then we conclude that $\|\rho_n(t, \cdot)\|_{L_x^2}^2$ converges in $C([0, T])$ by the Arzela-Ascoli theorem. Moreover, from (3.37), (3.39), (3.45)–(3.46), and Lemma 3.4, we have

$$\begin{aligned} \|\rho_n(t, \cdot)\|_{L_x^2} &\rightarrow \|\rho(t, \cdot)\|_{L_x^2} \quad \text{for any } t, \\ \|\nabla \rho_n\|_{L^2(\mathcal{O}_T)} &\rightarrow \|\nabla \rho\|_{L^2(\mathcal{O}_T)}. \end{aligned}$$

Since $\nabla \rho_n \rightharpoonup \nabla \rho$, it yields

$$\nabla \rho_n \rightarrow \nabla \rho \quad \text{in } L^2(\mathcal{O}_T).$$

Then we have

$$\nabla \rho_n \cdot \nabla u_n \rightarrow \nabla \rho \cdot \nabla u \quad \text{in } \mathcal{D}'(\mathcal{O}_T).$$

In addition, by (3.40)–(3.41), we have

$$\begin{aligned} &\int_0^t \int_{\mathcal{O}} \nabla \cdot (F(Q_n)I_3 - \nabla Q_n \odot \nabla Q_n + Q_n \Delta Q_n - \Delta Q_n Q_n + \sigma_* c_n^2 Q_n) \cdot \psi \, dx ds \\ &\rightarrow \int_0^t \int_{\mathcal{O}} \nabla \cdot (F(Q)I_3 - \nabla Q \odot \nabla Q + Q \Delta Q - \Delta Q Q + \sigma_* c^2 Q) \cdot \psi \, dx ds. \end{aligned}$$

Then we can pass to the limit in equation (3.18) as $n \rightarrow \infty$. We deduce that the limit function (c, ρ, u, Q) is a weak solution of problem (3.1)–(3.10). Finally, let us summarize the results in this section in the following:

Proposition 3.1. *Suppose $\beta > \max\{4, \gamma\}$. Then there exists a weak solution (c, ρ, u, Q) of problem (3.1)–(3.10) with the same estimates as in Lemma 3.3 and $Q \in S_0^3$ a.e. in \mathcal{O}_T . Moreover, the energy inequality (3.21) and estimates (3.25)–(3.35) hold for (c, ρ, u, Q) . Finally, we can find $r > 1$ such that $\rho_t, \Delta \rho \in L^r(\mathcal{O}_T)$, and equation (3.2) is satisfied a.e. in \mathcal{O}_T .*

4. THE VANISHING ARTIFICIAL VISCOSITY LIMIT

In this section, we let $\varepsilon \rightarrow 0$ in (3.1)–(3.4). We denote $(c_\varepsilon, \rho_\varepsilon, u_\varepsilon, Q_\varepsilon)$ the solution of problem (3.1)–(3.10), which we have obtained in Proposition 3.1. However, unlike the previous step, we do not have the higher integrability of the density as in Lemma 3.4. The boundedness of ρ_ε in $L^\infty(0, T; L^\gamma(\mathcal{O}) \cap L^\beta(\mathcal{O}))$ can guarantee only that ρ_ε^β converges to a Radon measure as $\varepsilon \rightarrow 0$, which is not easy to deal with. Thus, it is essential to obtain the strong compactness of ρ_ε in $L^1(\mathcal{O}_T)$. In order to do this, we need to fully use the uniform estimates (3.25)–(3.35) for $(c_\varepsilon, \rho_\varepsilon, u_\varepsilon, Q_\varepsilon)$, since the system includes some highly nonlinear coupling terms, such as the last three terms in (4.5) when we prove the higher integrability of the density and the convergence of the last four terms in (4.29) to (4.30) in the proof of the convergence of the effective viscous flux. Moreover, we introduce the useful operator \mathcal{B} related to the equation: $\operatorname{div} v = f$. See [5, 6, 19] for the construction and the proof of the following properties of the operator \mathcal{B} : For the problem

$$\operatorname{div} v = f, \quad v|_{\partial \mathcal{O}} = 0, \tag{4.1}$$

there exists a linear operator $\mathcal{B} = [B_1, B_2, B_3]$ with the following properties:

- (i) $\mathcal{B} : \{f \in L^p(\mathcal{O}) : \int_{\mathcal{O}} f \, dx = 0\} \mapsto (W_0^{1,p}(\mathcal{O}))^3$ is a bounded linear operator such that, for any $1 < p < \infty$,

$$\|\mathcal{B}[f]\|_{W_0^{1,p}(\mathcal{O})} \leq C(p) \|f\|_{L^p(\mathcal{O})}; \tag{4.2}$$

- (ii) $v = \mathcal{B}[f]$ is a solution of problem (4.1);
- (iii) If there is a vector function $\mathbf{g} \in (L^r(\mathcal{O}))^3$ with $\mathbf{g} \cdot \vec{n}|_{\partial\mathcal{O}} = 0$, then

$$\|\mathcal{B}[\operatorname{div} \mathbf{g}]\|_{L^r(\mathcal{O})} \leq C(p)\|\mathbf{g}\|_{L^r(\mathcal{O})}, \quad (4.3)$$

where $r \in (1, \infty)$ is arbitrary.

4.1. Estimates of the density independent of ϵ . We take the quantities:

$$\psi(t)\mathcal{B}[\rho_\epsilon - \bar{m}], \quad (4.4)$$

with $\psi \in \mathcal{D}(0, T)$, $0 \leq \psi \leq 1$, and $\bar{m} = \frac{1}{|\mathcal{O}|} \int_{\mathcal{O}} \rho_0(x) dx$, as test functions for (3.3). Note that \bar{m} is a constant such that this test function is well defined. We have the following result.

Lemma 4.1. *Assume that $(c_\epsilon, \rho_\epsilon, u_\epsilon, Q_\epsilon)$ is the solution of problem (3.1)–(3.10) constructed in Proposition 3.1. Then*

$$\|\rho_\epsilon\|_{L^{\gamma+1}(\mathcal{O}_T)}^{\gamma+1} + \delta\|\rho_\epsilon\|_{L^{\beta+1}(\mathcal{O}_T)}^{\beta+1} \leq C,$$

with C independent of ϵ .

Proof. The proof is similar to Lemma 3.1 in [15]. Let us apply the test function (4.4) to (3.3). Then, by a direct calculation, we have

$$\begin{aligned} & \int_0^T \int_{\mathcal{O}} \psi (\rho_\epsilon^{\gamma+1} + \delta \rho_\epsilon^{\beta+1}) dx dt \\ &= \bar{m} \int_0^T \psi \int_{\mathcal{O}} (\rho_\epsilon^\gamma + \delta \rho_\epsilon^\beta) dx dt + (\mu + \nu) \int_0^T \psi \int_{\mathcal{O}} (\rho_\epsilon - \bar{m}) \operatorname{div} u_\epsilon dx dt \\ & \quad - \int_0^T \psi_t \int_{\mathcal{O}} \rho_\epsilon u_\epsilon \cdot \mathcal{B}[\rho_\epsilon - \bar{m}] dx dt + \mu \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\epsilon^i \partial_j B_i[\rho_\epsilon - \bar{m}] dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} \rho_\epsilon u_\epsilon^i u_\epsilon^j \partial_j B_i[\rho_\epsilon - \bar{m}] dx dt - \epsilon \int_0^T \psi \int_{\mathcal{O}} \rho_\epsilon u_\epsilon \cdot \mathcal{B}(\Delta \rho_\epsilon) dx dt \\ & \quad + \int_0^T \psi \int_{\mathcal{O}} \rho_\epsilon u_\epsilon \cdot \mathcal{B}[\operatorname{div}(\rho_\epsilon u_\epsilon)] dx dt + \epsilon \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\epsilon^i \partial_j \rho_\epsilon B_i[\rho_\epsilon - \bar{m}] dx dt \\ & \quad + \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (\nabla Q_\epsilon \otimes \nabla Q_\epsilon - F(Q_\epsilon) \mathbf{I}_3) \cdot \mathcal{B}[\rho_\epsilon - \bar{m}] dx dt \\ & \quad + \sigma_* \int_0^T \psi \int_{\mathcal{O}} c_\epsilon^2 Q_\epsilon : \nabla \mathcal{B}[\rho_\epsilon - \bar{m}] dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (Q_\epsilon \Delta Q_\epsilon - \Delta Q_\epsilon Q_\epsilon) \cdot \mathcal{B}[\rho_\epsilon - \bar{m}] dx dt \\ &= \sum_{i=1}^{11} \mathcal{I}_i. \end{aligned} \quad (4.5)$$

Next, we estimate the terms on the right-hand side of the above equality by using the boundedness of solution $(c_\epsilon, \rho_\epsilon, u_\epsilon, Q_\epsilon)$ obtained in Proposition 3.1, in which the universal

constant $C > 0$ is independent of ε :

$$|\mathcal{I}_1| = \left| \bar{m} \int_0^T \psi \int_{\mathcal{O}} (\rho_\varepsilon^\gamma + \delta \rho_\varepsilon^\beta) dx dt \right| \leq C \|\psi\|_{L^\infty(0,T)} (\|\rho_\varepsilon\|_{L_t^\infty L_x^\gamma}^\gamma + \delta \|\rho_\varepsilon\|_{L_t^\infty L_x^\beta}^\beta) \leq C \|\psi\|_{L^\infty(0,T)},$$

$$\begin{aligned} |\mathcal{I}_2| &= (\mu + \nu) \left| \int_0^T \psi \int_{\mathcal{O}} (\rho_\varepsilon - \bar{m}) \operatorname{div} u_\varepsilon dx dt \right| \\ &\leq C \sqrt{T} \|\psi\|_{L^\infty(0,T)} (\|\rho_\varepsilon\|_{L_t^\infty L_x^2} + \bar{m} |\mathcal{O}|^{\frac{1}{2}}) \|\operatorname{div} u_\varepsilon\|_{L_{t,x}^2} \leq C \|\psi\|_{L^\infty(0,T)}, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_3| &= \left| \int_0^T \psi_t \int_{\mathcal{O}} \rho_\varepsilon u_\varepsilon \cdot \mathcal{B}[\rho_\varepsilon - \bar{m}] dx dt \right| \leq C \int_0^T |\psi_t| \|\sqrt{\rho_\varepsilon}\|_{L_x^4} \|\sqrt{\rho_\varepsilon} u_\varepsilon\|_{L_x^2} \|\mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^4} dt \\ &\leq C \|\rho_\varepsilon\|_{L_t^\infty L_x^2}^{\frac{1}{2}} \|\sqrt{\rho_\varepsilon} u_\varepsilon\|_{L_t^\infty L_x^2} \|\rho_\varepsilon - \bar{m}\|_{L_t^\infty L_x^2} \int_0^T |\psi_t| dt \leq C \|\psi_t\|_{L^1(0,T)}, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_4| &= \mu \left| \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\varepsilon^i \partial_j B_i [\rho_\varepsilon - \bar{m}] dx dt \right| \leq \mu \int_0^T |\psi| \|\nabla u_\varepsilon\|_{L_x^2} \|\nabla \mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^2} dt \\ &\leq \mu \int_0^T |\psi| \|\nabla u_\varepsilon\|_{L_x^2} \|\rho_\varepsilon - \bar{m}\|_{L_x^2} dt \leq C \|\psi\|_{L^\infty(0,T)}, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_5| &= \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\varepsilon u_\varepsilon^i u_\varepsilon^j \partial_j B_i [\rho_\varepsilon - \bar{m}] dx dt \right| \\ &\leq C \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6}^2 \|\nabla \mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^3} dt \\ &\leq C \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6}^2 \|\rho_\varepsilon - \bar{m}\|_{L_x^3} dt \leq C \|\psi\|_{L^\infty(0,T)}, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_6| &= \varepsilon \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\varepsilon u_\varepsilon \cdot \mathcal{B}[\Delta \rho_\varepsilon] dx dt \right| \leq \varepsilon \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6} \|\mathcal{B}[\Delta \rho_\varepsilon]\|_{L_x^2} dt \\ &\leq C \varepsilon \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6} \|\nabla \rho_\varepsilon\|_{L_x^2} dt \leq C \|\psi\|_{L^\infty(0,T)} \quad \text{for } \varepsilon < 1, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_7| &= \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\varepsilon u_\varepsilon \cdot \mathcal{B}[\operatorname{div}(\rho_\varepsilon u_\varepsilon)] dx dt \right| \leq \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6} \|\mathcal{B}[\operatorname{div}(\rho_\varepsilon u_\varepsilon)]\|_{L_x^2} dt \\ &\leq \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6} \|\rho_\varepsilon u_\varepsilon\|_{L_x^2} dt \leq C \int_0^T |\psi| \|\rho_\varepsilon\|_{L_x^3}^2 \|u_\varepsilon\|_{L_x^6}^2 dt \leq C \|\psi\|_{L^\infty(0,T)}. \end{aligned}$$

Since $\beta > 4$, by using the Sobolev embedding (Lemma A.3 in Appendix A), we have

$$\begin{aligned} \|B[\rho_\varepsilon - \bar{m}]\|_{L_x^\infty} &\leq C_1 \|\nabla B[\rho_\varepsilon - \bar{m}]\|_{L_x^\beta}^{\frac{3}{\beta}} \|B[\rho_\varepsilon - \bar{m}]\|_{L_x^\beta}^{1-\frac{3}{\beta}} + C_2 \|B[\rho_\varepsilon - \bar{m}]\|_{L_x^\beta} \\ &\leq C \|\rho_\varepsilon - \bar{m}\|_{L_x^\beta}. \end{aligned} \tag{4.6}$$

Then we have

$$\begin{aligned} |\mathcal{I}_8| &= \varepsilon \left| \int_0^T \psi \int_{\mathcal{O}} \nabla u_\varepsilon \cdot \mathcal{B}[\rho_\varepsilon - \bar{m}] \nabla \rho_\varepsilon dx dt \right| \\ &\leq \varepsilon \int_0^T |\psi| \|\nabla u_\varepsilon\|_{L_x^2} \|\nabla \rho_\varepsilon\|_{L_x^2} \|\mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^\infty} dt \end{aligned}$$

$$\leq \varepsilon \int_0^T |\psi| \|\nabla u_\varepsilon\|_{L_x^2} \|\nabla \rho_\varepsilon\|_{L_x^2} \|\rho_\varepsilon - \bar{m}\|_{L_x^\beta} dt \leq C \|\psi\|_{L^\infty(0,T)},$$

$$\begin{aligned} |\mathcal{I}_9| &= \left| \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (\nabla Q_\varepsilon \otimes \nabla Q_\varepsilon - F(Q_\varepsilon) I_3) \cdot \mathcal{B}[\rho_\varepsilon - \bar{m}] dx dt \right| \\ &\leq C \int_0^T |\psi| \left(\|\nabla Q_\varepsilon\|_{L_x^{\frac{10}{3}}}^2 \|\nabla \mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^{\frac{5}{2}}} + \left| \int_{\mathcal{O}} F(Q_\varepsilon) \operatorname{div} \mathcal{B}[\rho_\varepsilon - \bar{m}] dx \right| \right) dt \\ &\leq C \int_0^T |\psi| \left(\|\nabla Q_\varepsilon\|_{L_x^{\frac{10}{3}}}^2 \|\nabla \mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^{\frac{5}{2}}} + \left| \int_{\mathcal{O}} F(Q_\varepsilon) (\rho_\varepsilon - \bar{m}) dx \right| \right) dt \\ &\leq C \int_0^T |\psi| \left(\|\nabla Q_\varepsilon\|_{L_x^{\frac{10}{3}}}^2 \|\rho_\varepsilon - \bar{m}\|_{L_x^{\frac{5}{2}}} + (\|Q_\varepsilon\|_{L_x^5}^2 + \|Q_\varepsilon\|_{L_x^{10}}^4) \|\rho_\varepsilon - \bar{m}\|_{L_x^{\frac{5}{3}}} \right) dt \\ &\leq C \|\psi\|_{L^\infty(0,T)}, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_{10}| &= \sigma_* \left| \int_0^T \psi \int_{\mathcal{O}} c_\varepsilon^2 Q_\varepsilon \cdot \nabla \mathcal{B}[\rho_\varepsilon - \bar{m}] dx dt \right| \\ &\leq C \|c_\varepsilon\|_{L_{t,x}^\infty}^2 \int_0^T |\psi| \|Q_\varepsilon\|_{L_x^2} \|\nabla \mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^2} dt \leq C \|\psi\|_{L^\infty(0,T)}, \end{aligned}$$

$$\begin{aligned} |\mathcal{I}_{11}| &= \left| \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (Q_\varepsilon \Delta Q_\varepsilon - \Delta Q_\varepsilon Q_\varepsilon) \cdot \mathcal{B}[\rho_\varepsilon - \bar{m}] dx dt \right| \\ &\leq C \int_0^T |\psi| \|Q_\varepsilon\|_{L_x^4} \|\Delta Q_\varepsilon\|_{L_x^2} \|\nabla \mathcal{B}[\rho_\varepsilon - \bar{m}]\|_{L_x^4} dt \\ &\leq C \int_0^T |\psi| \|Q_\varepsilon\|_{L_x^4} \|\Delta Q_\varepsilon\|_{L_x^2} \|\rho_\varepsilon - \bar{m}\|_{L_x^4} dt \leq C \|\psi\|_{L^\infty(0,T)}. \end{aligned}$$

Combining all the above estimates together, we have

$$\int_0^T \int_{\mathcal{O}} \psi (\rho_\varepsilon^{\gamma+1} + \delta \rho_\varepsilon^{\beta+1}) dx dt \leq C (\|\psi\|_{L^\infty(0,T)} + \|\psi_t\|_{L^1(0,T)}). \quad (4.7)$$

Then we can take $\psi = \psi^k \in \mathcal{D}(0, T)$ in (4.7), with $\|\psi^k\|_{L^\infty(0,T)} + \|\psi_t^k\|_{L^1(0,T)} \leq C$ uniformly bounded independent of k , and $\psi^k(t) \rightarrow 1$ for $t \in (0, T)$ as $k \rightarrow \infty$, and then let $k \rightarrow \infty$ in the resulting (4.7) to conclude our desired result. \square

Remark 4.1. Lemma 4.1 implies that ρ_ε has higher integrability, which provides the weak convergence result for the pressure, i.e., $P_\varepsilon = \rho_\varepsilon^\gamma + \delta \rho_\varepsilon^\beta \rightharpoonup p$ in $L^{\frac{\beta+1}{\beta}}(\mathcal{O}_T)$.

4.2. Limit passage of $\varepsilon \rightarrow 0$. In this subsection, we fix parameter δ , and pass to the limit $\varepsilon \rightarrow 0$ in equations (3.1)–(3.4). To begin with, similarly to (3.40)–(3.41) in §3.3, we have

$$c_\varepsilon \rightharpoonup c \quad \text{in } L^2(0, T; H^1(\mathcal{O})), \quad c_\varepsilon \rightarrow c \quad \text{in } L^2(0, T; L^2(\mathcal{O})), \quad (4.8)$$

$$Q_\varepsilon \rightharpoonup Q \quad \text{in } L^2(0, T; H^2(\mathcal{O})), \quad Q_\varepsilon \rightarrow Q \quad \text{in } L^2(0, T; H^1(\mathcal{O})). \quad (4.9)$$

From the boundedness of ρ_ε in $L^{\beta+1}(\mathcal{O}_T)$, $\sqrt{\varepsilon}\nabla\rho_\varepsilon$ in $L^2(\mathcal{O}_T)$, and u_ε in $L^2(0, T; H_0^1(\mathcal{O}))$, we know that

$$\rho_\varepsilon \rightharpoonup \rho \quad \text{in } L^{\beta+1}(\mathcal{O}_T), \quad (4.10)$$

$$u_\varepsilon \rightharpoonup u \quad \text{in } L^2(0, T; H_0^1(\mathcal{O})), \quad (4.11)$$

$$\varepsilon \nabla \rho_\varepsilon \cdot \nabla u_\varepsilon \rightarrow 0 \quad \text{in } L^1(\mathcal{O}_T), \quad (4.12)$$

$$\varepsilon \Delta \rho_\varepsilon \rightarrow 0 \quad \text{in } L^2(0, T; H^{-1}(\mathcal{O})). \quad (4.13)$$

Moreover, we can also obtain the following convergence results as in §3.3:

$$\rho_\varepsilon \rightarrow \rho \quad \text{in } C([0, T]; L_{\text{weak}}^\beta(\mathcal{O})), \quad (4.14)$$

$$\rho_\varepsilon \rightarrow \rho \quad \text{in } C([0, T]; L_{\text{weak}}^\gamma(\mathcal{O})), \quad (4.15)$$

$$\rho_\varepsilon u_\varepsilon \rightarrow \rho u \quad \text{in } C([0, T]; L_{\text{weak}}^{\frac{2\gamma}{\gamma+1}}(\mathcal{O})), \quad (4.16)$$

$$\rho_\varepsilon u_\varepsilon \otimes u_\varepsilon \rightarrow \rho u \otimes u \quad \text{in } \mathcal{D}'(\mathcal{O}_T). \quad (4.17)$$

Finally, we conclude that the limit vector function (c, ρ, u, Q) satisfies the following equations in $\mathcal{D}'(\mathcal{O}_T)$:

$$\partial_t c + u \cdot \nabla c = D_0 c, \quad (4.18)$$

$$\partial_t \rho + \nabla \cdot (\rho u) = 0, \quad (4.19)$$

$$\partial_t (\rho u) + \nabla \cdot (\rho u \otimes u) + \nabla p = \mu \Delta u + (\nu + \mu) \nabla \operatorname{div} u + \nabla \cdot (F(Q) \mathbf{I}_3 - \nabla Q \odot \nabla Q) \quad (4.20)$$

$$+ \nabla \cdot (Q \Delta Q - \Delta Q Q) + \sigma_* \nabla \cdot (c^2 Q), \quad (4.21)$$

$$\partial_t Q + (u \cdot \nabla) Q + Q \Omega - \Omega Q = \Gamma H[Q, c], \quad (4.22)$$

along with the initial-boundary conditions (3.6)–(3.10), with

$$P_\varepsilon = \rho_\varepsilon^\gamma + \delta \rho_\varepsilon^\beta \rightharpoonup p \quad \text{in } L^{\frac{\beta+1}{\beta}}(\mathcal{O}_T) \quad (4.23)$$

for $\beta > \max\{4, \gamma\}$.

In the next step, we show that $p = \rho^\gamma + \delta \rho^\beta$, which is equivalent to the strong convergence of ρ_ε in $L^1(\mathcal{O}_T)$.

4.3. The effective viscous flux. The quantity, $\mathfrak{E}_\varepsilon := \rho_\varepsilon^\gamma + \delta \rho_\varepsilon^\beta - (\nu + 2\mu) \operatorname{div} u_\varepsilon$, is usually called the effective viscous flux, and its corresponding weak convergence limit is $\mathfrak{E} := p - (\nu + 2\mu) \operatorname{div} u$. The properties of \mathfrak{E}_ε (cf. [27, 33, 51]) play an important role in our problem. We introduce the following operator $\mathcal{A} = (\mathcal{A}_1, \mathcal{A}_2, \mathcal{A}_3) : \mathbb{R}^3 \rightarrow \mathbb{R}^3$:

$$\mathcal{A}_j[v] = \partial_j \Delta^{-1} v$$

with the Fourier transform:

$$\mathcal{F}(\mathcal{A}_j)(\xi) = -\frac{i\xi_j}{|\xi|^2},$$

and the following properties:

$$\operatorname{div} \mathcal{A}[v] = v, \quad \Delta \mathcal{A}_i[v] = \partial_i v, \quad (4.24)$$

$$\|\mathcal{A}_i[v]\|_{W^{1,p}(\mathcal{O})} \leq C(p, \mathcal{O}) \|v\|_{L^p(\mathbb{R}^3)} \quad \text{for } 1 < p < \infty, \quad (4.25)$$

$$\|\mathcal{A}_i[v]\|_{L^q(\mathcal{O})} \leq C(p, q, \mathcal{O})\|v\|_{L^p(\mathbb{R}^3)} \quad \text{if } p < \infty \text{ and } \frac{1}{q} \geq \frac{1}{p} - \frac{1}{3}, \quad (4.26)$$

$$\|\mathcal{A}_i[v]\|_{L^\infty(\mathcal{O})} \leq C(p, \mathcal{O})\|v\|_{L^p(\mathbb{R}^3)} \quad \text{for all } p > 3. \quad (4.27)$$

Lemma 4.2. *Let $(c_\varepsilon, \rho_\varepsilon, u_\varepsilon, Q_\varepsilon)$ be a sequence of solutions constructed in Proposition 3.1, and let (c, ρ, u, Q, p) be the limits satisfying (4.18)–(4.23). Then*

$$\lim_{\varepsilon \rightarrow 0^+} \int_0^T \psi \int_{\mathcal{O}} \phi \mathfrak{E}_\varepsilon \rho_\varepsilon \, dx dt = \int_0^T \psi \int_{\mathcal{O}} \phi \mathfrak{E} \rho \, dx dt,$$

for any $\psi \in \mathcal{D}(0, T)$ and $\phi \in \mathcal{D}(\mathcal{O})$.

Proof. We prove this lemma based on the div-curl lemma of compensated compactness. By Proposition 3.1, we know that $(\rho_\varepsilon, u_\varepsilon)$ satisfies (3.2) a.e. on \mathcal{O}_T with boundary condition (3.9). If $(\rho_\varepsilon, u_\varepsilon)$ are extended to be zero outside \mathcal{O} , then it satisfies

$$\partial_t \rho_\varepsilon + \nabla \cdot (\rho_\varepsilon u_\varepsilon) = \varepsilon \operatorname{div} (1_{\mathcal{O}} \nabla \rho_\varepsilon) \quad \text{in } \mathcal{D}'((0, T) \times \mathbb{R}^3),$$

with $1_{\mathcal{O}}$ the characteristic function on \mathcal{O} . Consider the following test function to (3.3) as

$$\varphi_\varepsilon(t, x) = \psi(t) \phi(x) \mathcal{A}[\rho_\varepsilon], \quad (4.28)$$

with $\psi \in \mathcal{D}(0, T)$ and $\phi \in \mathcal{D}(\mathcal{O})$. Similarly to (4.5), we apply the test function (4.28) to (3.3). By a direct calculation, we have

$$\begin{aligned} & \int_0^T \psi \int_{\mathcal{O}} \phi (\rho_\varepsilon^{\gamma+1} + \delta \rho_\varepsilon^{\beta+1} - (\nu + 2\mu) \operatorname{div} u_\varepsilon) \rho_\varepsilon \, dx dt \\ &= - \int_0^T \psi \int_{\mathcal{O}} (\rho_\varepsilon^\gamma + \delta \rho_\varepsilon^\beta) \partial_i \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt + (\mu + \nu) \int_0^T \psi \int_{\mathcal{O}} \operatorname{div} u_\varepsilon \partial_i \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt \\ & \quad + \mu \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\varepsilon^i \partial_j \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt - \mu \int_0^T \psi \int_{\mathcal{O}} u_\varepsilon^i \partial_j \phi \partial_j \mathcal{A}_i[\rho_\varepsilon] \, dx dt \\ & \quad + \mu \int_0^T \psi \int_{\mathcal{O}} u_\varepsilon^i \partial_i \phi \rho_\varepsilon \, dx dt - \varepsilon \int_0^T \psi \int_{\mathcal{O}} \phi \rho_\varepsilon u_\varepsilon^i \mathcal{A}_i \operatorname{div} (1_{\mathcal{O}} \nabla \rho_\varepsilon) \, dx dt \\ & \quad + \int_0^T \psi \int_{\mathcal{O}} \phi \rho_\varepsilon u_\varepsilon^i (\mathcal{A}_i[\operatorname{div}(\rho_\varepsilon u_\varepsilon)] - u_\varepsilon^j \partial_j \mathcal{A}_i[\rho_\varepsilon]) \, dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} \phi \rho_\varepsilon u_\varepsilon^i \mathcal{A}_i[\rho_\varepsilon] \, dx dt - \int_0^T \psi \int_{\mathcal{O}} \rho_\varepsilon u_\varepsilon^i u_\varepsilon^j \partial_j \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt \\ & \quad + \varepsilon \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\varepsilon^i \partial_j \rho_\varepsilon \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt + \sigma_* \int_0^T \psi \int_{\mathcal{O}} c_\varepsilon^2 Q_\varepsilon^{ij} (\partial_j \phi \mathcal{A}_i[\rho_\varepsilon] + \phi \partial_j \mathcal{A}_i[\rho_\varepsilon]) \, dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} (\partial_i Q_\varepsilon^{kl} \partial_j Q_\varepsilon^{kl} - F(Q_\varepsilon) \delta_{ij}) \partial_j \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} (\partial_i Q_\varepsilon^{kl} \partial_j Q_\varepsilon^{kl} - F(Q_\varepsilon) \delta_{ij}) \phi \partial_j \mathcal{A}_i[\rho_\varepsilon] \, dx dt \\ & \quad + \int_0^T \psi \int_{\mathcal{O}} (Q_\varepsilon^{ik} \Delta Q_\varepsilon^{kj} - \Delta Q_\varepsilon^{ik} Q_\varepsilon^{kj}) (\partial_j \phi \mathcal{A}_i[\rho_\varepsilon] + \phi \partial_j \mathcal{A}_i[\rho_\varepsilon]) \, dx dt \\ &= \sum_{i=1}^{14} \mathcal{I}_i. \end{aligned} \quad (4.29)$$

Remark 4.2. The function ρ_ε extended by zero outside \mathcal{O} admits the time derivative as

$$\partial_t \rho_\varepsilon = \begin{cases} \varepsilon \Delta \rho_\varepsilon - \operatorname{div}(\rho_\varepsilon u_\varepsilon) & \text{in } \mathcal{O}, \\ 0 & \text{in } \mathbb{R}^3 \setminus \mathcal{O}. \end{cases}$$

Since $u_\varepsilon|_{\partial\mathcal{O}} = 0$, we have

$$\operatorname{div}(\rho_\varepsilon u_\varepsilon) = 0 \quad \text{on } \mathbb{R}^3 \setminus \mathcal{O}.$$

Moreover, since $\nabla \rho_\varepsilon \cdot \vec{n}|_{\partial\mathcal{O}} = 0$ and $\Delta \rho_\varepsilon \in L^r(\mathcal{O}_T)$ for some $r > 1$, we have

$$\operatorname{div}(1_{\mathcal{O}} \nabla \rho_\varepsilon) = \begin{cases} \Delta \rho_\varepsilon & \text{in } \mathcal{O}, \\ 0 & \text{in } \mathbb{R}^3 \setminus \mathcal{O}. \end{cases}$$

Next, we do the zero extension to the limit function ρ to \mathbb{R}^3 , and repeat the same procedure above. This step is guaranteed by the following results from [15]:

Lemma 4.3. *Assume that $(\rho, u) \in L^2(\mathcal{O}_T) \times L^2(0, T; H_0^1(\mathcal{O}))$ is a solution of (4.19) in $\mathcal{D}'(\mathcal{O}_T)$. Then, extending (ρ, u) to be zero in $\mathbb{R}^3 \setminus \mathcal{O}$, equation (4.19) still holds in $\mathcal{D}'((0, T) \times \mathbb{R}^3)$.*

Here, let us apply the test function $\varphi = \psi(t)\phi(x)\mathcal{A}[\rho]$ to (4.21). By a similar calculation as before, we have

$$\begin{aligned} & \int_0^T \psi \int_{\mathcal{O}} \phi(p - (\nu + 2\mu)\operatorname{div} u) \rho \, dx dt \\ &= - \int_0^T \psi \int_{\mathcal{O}} p \partial_i \phi \mathcal{A}_i[\rho] \, dx dt + (\mu + \nu) \int_0^T \psi \int_{\mathcal{O}} \operatorname{div} u \partial_i \phi \mathcal{A}_i[\rho] \, dx dt \\ & \quad + \mu \int_0^T \psi \int_{\mathcal{O}} \partial_j u^i \partial_j \phi \mathcal{A}_i[\rho] \, dx dt - \mu \int_0^T \psi \int_{\mathcal{O}} u^i \partial_j \phi \partial_j \mathcal{A}_i[\rho] \, dx dt \\ & \quad + \mu \int_0^T \psi \int_{\mathcal{O}} u^i \partial_i \phi \rho \, dx dt + \int_0^T \psi \int_{\mathcal{O}} \phi \rho u^i (\mathcal{A}_i[\operatorname{div}(\rho u)] - u^j \partial_j \mathcal{A}_i[\rho]) \, dx dt \\ & \quad - \int_0^T \psi_t \int_{\mathcal{O}} \phi \rho u^i \mathcal{A}_i[\rho] \, dx dt - \int_0^T \psi \int_{\mathcal{O}} \rho u^i u^j \partial_j \phi \mathcal{A}_i[\rho] \, dx dt \\ & \quad + \sigma_* \int_0^T \psi \int_{\mathcal{O}} c^2 Q^{ij} (\partial_j \phi \mathcal{A}_i[\rho] + \phi \partial_j \mathcal{A}_i[\rho]) \, dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} (\partial_i Q^{kl} \partial_j Q^{kl} - F(Q) \delta_{ij}) \partial_j \phi \mathcal{A}_i[\rho] \, dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} (\partial_i Q^{kl} \partial_j Q^{kl} - F(Q) \delta_{ij}) \phi \partial_j \mathcal{A}_i[\rho] \, dx dt \\ & \quad + \int_0^T \psi \int_{\mathcal{O}} (Q^{ik} \Delta Q^{kj} - \Delta Q^{ik} Q^{kj}) (\partial_j \phi \mathcal{A}_i[\rho] + \phi \partial_j \mathcal{A}_i[\rho]) \, dx dt \\ &= \sum_{j=1}^{12} \mathcal{J}_j. \end{aligned} \tag{4.30}$$

Next, in order to prove Lemma 4.2, we need to show that the right-hand side of (4.29) converges to the right-hand side of (4.30). First, by the uniform bound of ρ_ε in $L^\infty(0, T; L^\beta(\mathcal{O}))$, similarly to (3.43), we have

$$\rho_\varepsilon \rightarrow \rho \quad \text{in } C([0, T]; L_{\text{weak}}^\beta(\mathcal{O})) \quad \text{as } \varepsilon \rightarrow 0.$$

Then, for any $\phi \in (L^\beta(\mathcal{O}))^*$, we extend ϕ by zero outside of \mathcal{O} to obtain

$$\begin{aligned} (\partial_j \mathcal{A}_i[\rho_\varepsilon], \phi)_{L^\beta(\mathcal{O}) \times (L^\beta(\mathcal{O}))^*} &= \int_{\mathbb{R}^3} \partial_j \mathcal{A}_i[\rho_\varepsilon] \phi \, dx = \int_{\mathbb{R}^3} \frac{\hat{\rho}_\varepsilon \xi_i \xi_j}{|\xi|^2} \hat{\phi} \, d\xi \\ &= \int_{\mathbb{R}^3} \rho_\varepsilon \partial_j \mathcal{A}_i[\phi] \, dx = (\rho_\varepsilon, \partial_j \mathcal{A}_i[\phi])_{L^\beta(\mathcal{O}) \times (L^\beta(\mathcal{O}))^*} \\ &\longrightarrow (\rho, \partial_j \mathcal{A}_i[\phi])_{L^\beta(\mathcal{O}) \times (L^\beta(\mathcal{O}))^*} = (\partial_j \mathcal{A}_i[\rho], \phi)_{L^\beta(\mathcal{O}) \times (L^\beta(\mathcal{O}))^*} \quad \text{as } \varepsilon \rightarrow 0, \end{aligned}$$

from which we infer

$$\mathcal{A}[\rho_\varepsilon] \rightarrow \mathcal{A}[\rho] \quad \text{in } C([0, T], W_{\text{weak}}^{1, \beta}(\mathcal{O})).$$

This, combining with Proposition A.1 ($\beta > \frac{6}{5}$) and the compact imbedding $W^{1, \beta}(\mathcal{O}) \Subset C(\bar{\mathcal{O}})$, gives

$$\nabla \mathcal{A}[\rho_\varepsilon] \rightarrow \nabla \mathcal{A}[\rho] \quad \text{in } C([0, T]; H^{-1}(\mathcal{O})), \quad (4.31)$$

$$\mathcal{A}[\rho_\varepsilon] \rightarrow \mathcal{A}[\rho] \quad \text{in } C(\bar{\mathcal{O}}_T). \quad (4.32)$$

From (4.11) and (4.32), we see that, as $\varepsilon \rightarrow 0$,

$$\mathcal{I}_2 \rightarrow \mathcal{J}_2, \quad \mathcal{I}_3 \rightarrow \mathcal{J}_3.$$

By (4.16) and (4.32), we find that, as $\varepsilon \rightarrow 0$,

$$\mathcal{I}_8 \rightarrow \mathcal{J}_7.$$

From (4.11) and (4.16), we have

$$\|\rho_\varepsilon u_\varepsilon \otimes u_\varepsilon\|_{L_x^{\frac{6\gamma}{3+4\gamma}}} \leq C \|\rho_\varepsilon u_\varepsilon\|_{L_x^{\frac{2\gamma}{\gamma+1}}} \|u_\varepsilon\|_{L_x^6},$$

which, combined with (4.17), gives

$$\rho_\varepsilon u_\varepsilon \otimes u_\varepsilon \rightharpoonup \rho u \otimes u \quad \text{in } L^2(0, T; L^{\frac{6\gamma}{3+4\gamma}}(\mathcal{O})).$$

Thus, together with (4.32), we see that, if $\beta > \frac{6\gamma}{2\gamma-3}$,

$$\mathcal{I}_9 \rightarrow \mathcal{J}_8 \quad \text{as } \varepsilon \rightarrow 0.$$

Combining (4.11) with (4.31), we obtain that, as $\varepsilon \rightarrow 0$,

$$\mathcal{I}_4 \rightarrow \mathcal{J}_4.$$

In the same way as above, since $\beta > 4$, we know

$$\mathcal{I}_5 \rightarrow \mathcal{J}_5.$$

It follows from the boundedness of ρ_ε in $L^\beta(\mathcal{O})$, $\rho_\varepsilon u_\varepsilon$ in $L^{\frac{2\gamma}{\gamma+1}}(\mathcal{O})$, and (4.25) that

$$\begin{aligned} &\|\rho_\varepsilon \partial_j \mathcal{A}_i[\rho_\varepsilon u_\varepsilon^j] - \rho_\varepsilon u_\varepsilon^j \partial_j \mathcal{A}_i[\rho_\varepsilon]\|_{L_x^\alpha} \\ &\leq \|\rho_\varepsilon\|_{L_x^\beta} \|\partial_j \mathcal{A}_i[\rho_\varepsilon u_\varepsilon^j]\|_{L_x^{\frac{2\gamma}{\gamma+1}}} + \|\rho_\varepsilon u_\varepsilon\|_{L_x^{\frac{2\gamma}{\gamma+1}}} \|\partial_j \mathcal{A}_i[\rho_\varepsilon]\|_{L_x^\beta} \end{aligned}$$

$$\leq \|\rho_\varepsilon\|_{L_x^\beta} \|\rho_\varepsilon u_\varepsilon\|_{L_x^{\frac{2\gamma}{\gamma+1}}} \leq C,$$

where $\frac{1}{\alpha} = \frac{\gamma+1}{2\gamma} + \frac{1}{\beta} < \frac{5}{6}$, if $\beta > \frac{6\gamma}{2\gamma-3}$. Then we conclude

$$\rho_\varepsilon \partial_j \mathcal{A}_i[\rho_\varepsilon u_\varepsilon^j] - \rho_\varepsilon u_\varepsilon^j \partial_j \mathcal{A}_i[\rho_\varepsilon] \in L^\infty(0, T; L^\alpha(\mathcal{O})).$$

From Lemma 3.4 in [15] and the compact embedding of $L^\alpha(\mathcal{O})$ in $H^{-1}(\mathcal{O})$, we infer that

$$\rho_\varepsilon \partial_j \mathcal{A}_i[\rho_\varepsilon u_\varepsilon^j] - \rho_\varepsilon u_\varepsilon^j \partial_j \mathcal{A}_i[\rho_\varepsilon] \rightarrow \rho \partial_j \mathcal{A}_i[\rho u^j] - \rho u^j \partial_j \mathcal{A}_i[\rho] \quad \text{in } H^{-1}(\mathcal{O}).$$

Then, after applying Lebesgue convergence theorem, we obtain

$$\rho_\varepsilon \partial_j \mathcal{A}_i[\rho_\varepsilon u_\varepsilon^j] - \rho_\varepsilon u_\varepsilon^j \partial_j \mathcal{A}_i[\rho_\varepsilon] \rightarrow \rho \partial_j \mathcal{A}_i[\rho u^j] - \rho u^j \partial_j \mathcal{A}_i[\rho] \quad \text{in } L^2(0, T; H^{-1}(\mathcal{O})),$$

which, combined with (4.11), yields

$$\mathcal{I}_7 \rightarrow \mathcal{J}_6 \quad \text{as } \varepsilon \rightarrow 0.$$

Moreover, we have

$$\begin{aligned} |\mathcal{I}_6| &= \varepsilon \left| \int_0^T \psi \int_{\mathcal{O}} \phi \rho_\varepsilon u_\varepsilon^i \mathcal{A}_i[\operatorname{div}(1_\Omega \nabla \rho_\varepsilon)] \, dx dt \right| \\ &\leq C\varepsilon \int_0^T \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6} \|\mathcal{A}_i[\operatorname{div}(1_\Omega \nabla \rho_\varepsilon)]\|_{L_x^2} \, dt \\ &\leq C\varepsilon \int_0^T \|\rho_\varepsilon\|_{L_x^3} \|u_\varepsilon\|_{L_x^6} \|\nabla \rho_\varepsilon\|_{L_x^2} \, dt \\ &\leq C\varepsilon \|\rho_\varepsilon\|_{L_t^\infty L_x^3} \|u_\varepsilon\|_{L_t^2 L_x^6} \|\nabla \rho_\varepsilon\|_{L_t^2 L_x^2} \leq C\sqrt{\varepsilon} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0. \end{aligned}$$

Since $\beta > 3$, we have

$$\begin{aligned} |\mathcal{I}_{10}| &= \varepsilon \left| \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\varepsilon^i \partial_j \rho_\varepsilon \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt \right| \\ &\leq C\varepsilon \int_0^T \|\nabla u_\varepsilon\|_{L_x^2} \|\nabla \rho_\varepsilon\|_{L_x^2} \|\mathcal{A}_i[\rho_\varepsilon]\|_{L_x^\infty} \, dt \\ &\leq C\varepsilon \int_0^T \|\nabla u_\varepsilon\|_{L_x^2} \|\nabla \rho_\varepsilon\|_{L_x^2} \|\rho_\varepsilon\|_{L_x^\beta} \, dt \\ &\leq C\varepsilon \|\nabla u_\varepsilon\|_{L_t^2 L_x^2} \|\nabla \rho_\varepsilon\|_{L_t^2 L_x^2} \|\rho_\varepsilon\|_{L_t^\infty L_x^\beta} \leq C\sqrt{\varepsilon} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0. \end{aligned}$$

From (4.8)–(4.9) and (4.31)–(4.32), we have

$$\mathcal{I}_{11} = \sigma_* \int_0^T \psi \int_{\mathcal{O}} c_\varepsilon^2 Q_\varepsilon^{ij} (\partial_j \phi \mathcal{A}_i[\rho_\varepsilon] + \phi \partial_j \mathcal{A}_i[\rho_\varepsilon]) \, dx dt \rightarrow \mathcal{J}_9.$$

From (4.9), (4.32), and the boundedness of Q_ε in $L^2(0, T; H^2(\mathcal{O}))$, we have

$$\mathcal{I}_{12} = \int_0^T \psi \int_{\mathcal{O}} (\partial_i Q_\varepsilon^{kl} \partial_j Q_\varepsilon^{kl} - F(Q_\varepsilon) \delta_{ij}) \partial_j \phi \mathcal{A}_i[\rho_\varepsilon] \, dx dt \rightarrow \mathcal{J}_{10},$$

$$\mathcal{I}_{13} = \int_0^T \psi \int_{\mathcal{O}} (\partial_i Q_\varepsilon^{kl} \partial_j Q_\varepsilon^{kl} - F(Q_\varepsilon) \delta_{ij}) \phi \partial_j \mathcal{A}_i[\rho_\varepsilon] \, dx dt \rightarrow \mathcal{J}_{11},$$

and

$$\mathcal{I}_{14} = - \int_0^T \psi \int_{\mathcal{O}} (Q_\varepsilon^{ik} \Delta Q_\varepsilon^{kj} - \Delta Q_\varepsilon^{ik} Q_\varepsilon^{kj}) (\partial_j \phi \mathcal{A}_i[\rho_\varepsilon] + \phi \partial_j \mathcal{A}_i[\rho_\varepsilon]) \, dx dt \rightarrow \mathcal{I}_{12}.$$

Next, following the same argument as in Subsection 3.5 in [14], we obtain the strong convergence of the density sequence ρ_ε in $L^1(\mathcal{O}_T)$, *i.e.*

$$p = \rho^\gamma + \delta \rho^\beta. \quad (4.33)$$

□

Then we have the following proposition:

Proposition 4.1. *Let $\beta > \max\{\frac{6\gamma}{2\gamma-3}, \gamma, 4\}$. Then, for any given $T > 0$ and $\delta > 0$, there exists a finite-energy weak solution (c, ρ, u, Q) of the problem:*

$$\partial_t c + u \cdot \nabla c = D_0 \Delta c, \quad (4.34)$$

$$\partial_t \rho + \operatorname{div}(\rho u) = 0, \quad (4.35)$$

$$\begin{aligned} \partial_t(\rho u) + \nabla \cdot (\rho u \otimes u) + \nabla(\rho^\gamma + \delta \rho^\beta) &= \mu \Delta u + (\nu + \mu) \nabla \operatorname{div} u + \nabla \cdot (Q \Delta Q - \Delta Q Q) \\ &\quad + \nabla \cdot (F(Q) \mathbf{I}_3 - \nabla Q \odot \nabla Q) + \sigma_* \nabla \cdot (c^2 Q), \end{aligned} \quad (4.36)$$

$$\partial_t Q + (u \cdot \nabla) Q + Q \Omega - \Omega Q = \Gamma H[Q, c], \quad (4.37)$$

with initial-boundary conditions (3.5)–(3.10). Moreover, equation (4.35) holds in the sense of renormalized solutions on $\mathcal{D}'((0, T) \times \mathbb{R}^3)$, provided that (ρ, u) are extended to be zero on $\mathbb{R}^3 \setminus \mathcal{O}$. In addition, the following estimates are valid:

$$0 < \underline{c} \leq c(x, t) \leq \bar{c}, \quad \|c\|_{L^\infty(0, T; L^2(\mathcal{O})) \cap L^2(0, T; H^1(\mathcal{O}))} \leq C, \quad (4.38)$$

$$\sup_{t \in [0, T]} \|\rho(t, \cdot)\|_{L^\gamma(\mathcal{O})}^\gamma \leq C, \quad (4.39)$$

$$\delta \sup_{t \in [0, T]} \|\rho(t, \cdot)\|_{L^\beta(\mathcal{O})}^\beta \leq C, \quad (4.40)$$

$$\sup_{t \in [0, T]} \|\sqrt{\rho}(t, \cdot) u(t, \cdot)\|_{L^2(\mathcal{O})}^2 \leq C, \quad (4.41)$$

$$\|u\|_{L^2(0, T; H_0^1(\mathcal{O}))} \leq C, \quad (4.42)$$

$$\|Q\|_{L^\infty(0, T; H^1(\mathcal{O}) \cap L^4(\mathcal{O})) \cap L^2(0, T; H^2(\mathcal{O}) \cap L^6(\mathcal{O}))} \leq C, \quad (4.43)$$

$$\|Q\|_{L^{10}(\mathcal{O}_T)} \leq C, \quad (4.44)$$

$$\|\nabla Q\|_{L^{\frac{10}{3}}(\mathcal{O}_T)} \leq C, \quad (4.45)$$

where C is a constant independent of ε .

Remark 4.3. The initial conditions (3.6)–(3.8) are satisfied in the weak sense, since

$$\rho_\varepsilon \rightarrow \rho \text{ in } C([0, T]; L_{\text{weak}}^\beta(\mathcal{O})), \quad \rho_\varepsilon u_\varepsilon \rightarrow \rho u \text{ in } C([0, T]; L_{\text{weak}}^{\frac{2\gamma}{\gamma+1}}(\mathcal{O})),$$

from (4.10)–(4.11).

5. PASSING TO THE LIMIT IN THE ARTIFICIAL PRESSURE

We denote by $(c_\delta, \rho_\delta, u_\delta, Q_\delta)$ the approximate solutions constructed in Proposition 4.1. In this section, we let $\delta \rightarrow 0$ in (4.34)–(4.37) to obtain the solution of the original problem (1.2)–(1.5).

In order for the solution (ρ_δ, u_δ) to satisfy the initial condition (3.6)–(3.7) in Proposition 4.1, we first modify the general initial data (ρ_0, m_0) to satisfy the compatibility condition (1.8). As in [15], it is easy to find a sequence $\tilde{\rho}_\delta \in C_0^3(\mathcal{O})$ with the property:

$$0 \leq \tilde{\rho}_\delta(x) \leq \frac{1}{2}\delta^{-\frac{1}{\beta}}, \quad \|\tilde{\rho}_\delta - \rho_0\|_{L^2(\mathcal{O})} < \delta.$$

Take $\rho_{0,\delta} = \tilde{\rho}_\delta + \delta$. From (3.6)–(3.7), we have

$$0 < \delta \leq \rho_{0,\delta}(x) \leq \delta^{-\frac{1}{\beta}}, \quad \nabla \rho_{0,\delta} \cdot \vec{n}|_{\partial\mathcal{O}} = 0, \quad (5.1)$$

with

$$\rho_{0,\delta} \rightarrow \rho_0 \quad \text{in } L^\gamma(\mathcal{O}) \text{ as } \delta \rightarrow 0. \quad (5.2)$$

Set

$$\tilde{q}_\delta(x) = \begin{cases} m_0(x) \sqrt{\frac{\rho_{0,\delta}}{\rho_0}} & \text{if } \rho_0(x) > 0, \\ 0 & \text{if } \rho_0(x) = 0. \end{cases}$$

Then it follows from (1.8) that $\frac{|\tilde{q}_\delta|^2}{\rho_{0,\delta}}$ is uniformly bounded in $L^1(\mathcal{O})$. It is also direct to find $h_\delta \in C^2(\bar{\mathcal{O}})$ such that

$$\left\| \frac{\tilde{q}_\delta}{\sqrt{\rho_{0,\delta}}} - h_\delta \right\|_{L^2(\mathcal{O})} < \delta.$$

Taking $m_{0,\delta} = h_\delta \sqrt{\rho_{0,\delta}}$, we can check that

$$\frac{|m_{0,\delta}|^2}{\rho_{0,\delta}} \quad \text{is uniformly bounded in } L^1(\mathcal{O}) \text{ with respect to } \delta > 0, \quad (5.3)$$

and

$$m_{0,\delta} \rightarrow m_0 \quad \text{in } L^1(\mathcal{O}) \text{ as } \delta \rightarrow 0. \quad (5.4)$$

From now on, we deal with the sequence of approximate solutions $(c_\delta, \rho_\delta, u_\delta, Q_\delta)$ of problem (4.34)–(4.37) with the initial data $(c_0, \rho_{0,\delta}, m_{0,\delta}, Q_0)$. The existence of such a solution is provided by Proposition 4.1. We notice that all the corresponding estimates in Proposition 4.1 are independent of δ , by virtue of (5.1) and (5.3).

5.1. On the integrability of the density. Now we provide some pressure estimates independent of $\delta > 0$. From (4.39), $\rho_\delta \in L^\infty(0, T; L^\gamma(\Omega))$ uniformly in δ yields that $\rho_\delta^\gamma \rightharpoonup \mu$ by measure, as $\delta \rightarrow 0$. Hence, we need the higher integrability of ρ_δ . The technique is similar to that in §4.1. In addition, we need to make estimates carefully for the highly nonlinear terms like the last three terms in (5.7)–(5.8) by using the estimates of the solution obtained in Proposition 4.1.

Since the continuity equation (4.35) is satisfied in the sense of renormalized solutions in $\mathcal{D}'((0, T) \times \mathbb{R}^3)$, we can apply the standard mollifying operator on both sides of (4.35) to obtain

$$\partial_t S_m[g(\rho_\delta)] + \operatorname{div}(S_m[g(\rho_\delta)u_\delta]) + S_m[(g'(\rho_\delta)\rho_\delta - g(\rho_\delta))\operatorname{div} u_\delta] = r_m,$$

with

$$r_m \rightarrow 0 \quad \text{in } L^2(0, T; L^2(\mathbb{R}^3)) \text{ as } m \rightarrow \infty.$$

As in §4.1, we use operator \mathcal{B} to construct the test function as

$$\phi(t, x) = \psi(t) \mathcal{B}[S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy], \quad (5.5)$$

with

$$\psi \in \mathcal{D}(0, T), \quad \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy := \frac{1}{|\mathcal{O}|} \int_{\mathcal{O}} S_m[g(\rho_\delta)] dy.$$

Since $\phi|_{\partial\mathcal{O}} = 0$, $\phi \in L^\infty(0, T; H_0^1(\mathcal{O}))$, and $\partial_t \phi \in L^2(0, T; H_0^1(\mathcal{O}))$, we know that ϕ can be used as a test function for (4.36). We approximate the function: $g(z) = z^\theta$, $0 < \theta \leq 1$, by a sequence of functions $\{z^\theta \chi_n(z)\}_{n=1}^\infty$, where $\{\chi_n(z)\}_{n=1}^\infty$ are cutoff functions with

$$\chi_n(z) = \begin{cases} 1 & \text{if } z \in [0, n], \\ 0 & \text{if } z > 2n. \end{cases}$$

Then, from (4.38)–(4.45), we employ the same techniques as in Lemma 4.1 to obtain the following result:

Lemma 5.1. *If $\gamma > \frac{3}{2}$, there exists a constant $\theta \in (0, \min\{\frac{1}{4}, \frac{2\gamma}{3} - 1\})$ depending only on γ , and $C > 0$ independent of δ such that*

$$\int_0^T \int_{\mathcal{O}} (\rho_\delta^{\gamma+\theta} + \delta \rho_\delta^{\beta+\theta}) dx dt \leq C. \quad (5.6)$$

Proof. Applying the test function (5.5) with the choice of $\psi = \psi(t) \geq 0$ to (4.36), we have

$$\begin{aligned} & \int_0^T \psi(t) \int_{\mathcal{O}} (\rho_\delta^\gamma + \delta \rho_\delta^\beta) S_m[g(\rho_\delta)] dx dt \\ &= \int_0^T \psi(t) \int_{\mathcal{O}} (\rho_\delta^\gamma + \delta \rho_\delta^\beta) dx \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy dt + (\mu + \nu) \int_0^T \psi \int_{\mathcal{O}} \operatorname{div} u_\delta S_m[g(\rho_\delta)] dx dt \\ & \quad - \int_0^T \psi_t \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\ & \quad + \mu \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\delta^i \partial_j B_i [S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta^i u_\delta^j \partial_j B_i [S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[r_m - \oint_{\mathcal{O}} r_m dy] dx dt \\ & \quad - \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[S_m[(g'(\rho_\delta)\rho_\delta - g(\rho_\delta)) \operatorname{div} u_\delta] \\ & \quad \quad - \oint_{\mathcal{O}} S_m[(g'(\rho_\delta)\rho_\delta - g(\rho_\delta)) \operatorname{div} u_\delta] dy] dx dt \\ & \quad + \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[\operatorname{div}(S_m[g(\rho_\delta)] u_\delta)] dx dt \end{aligned}$$

$$\begin{aligned}
& - \int_0^T \psi \int_{\mathcal{O}} (\nabla Q_\delta \otimes \nabla Q_\delta - F(Q_\delta) \mathbf{I}_3) : \nabla \mathcal{B}[S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\
& + \sigma_* \int_0^T \psi \int_{\mathcal{O}} c_\delta^2 Q_\delta : \nabla \mathcal{B}[S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\
& - \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (Q_\delta \Delta Q_\delta - \Delta Q_\delta Q_\delta) \cdot \mathcal{B}[S_m[g(\rho_\delta)] - \oint_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\
& = \sum_{i=1}^{11} \mathcal{I}_i.
\end{aligned} \tag{5.7}$$

By the estimates in Proposition 4.1, we can pass to the limit in (5.7) as $m \rightarrow \infty$ to obtain

$$\begin{aligned}
& \int_0^T \psi(t) \int_{\mathcal{O}} (\rho_\delta^\gamma + \delta \rho_\delta^\beta) g(\rho_\delta) dx dt \\
& = \int_0^T \psi(t) \int_{\mathcal{O}} (\rho_\delta^\gamma + \delta \rho_\delta^\beta) dx \oint_{\mathcal{O}} g(\rho_\delta) dy dt + (\mu + \nu) \int_0^T \psi \int_{\mathcal{O}} \operatorname{div} u_\delta g(\rho_\delta) dx dt \\
& \quad - \int_0^T \psi_t \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[g(\rho_\delta) - \oint_{\mathcal{O}} g(\rho_\delta) dy] dx dt \\
& \quad + \mu \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\delta^i \partial_j B_i [g(\rho_\delta) - \oint_{\mathcal{O}} g(\rho_\delta) dy] dx dt \\
& \quad - \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta^i \partial_j^i \partial_j B_i [g(\rho_\delta) - \oint_{\mathcal{O}} g(\rho_\delta) dy] dx dt \\
& \quad - \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[(g'(\rho_\delta) \rho_\delta - g(\rho_\delta)) \operatorname{div} u_\delta - \oint_{\mathcal{O}} (g'(\rho_\delta) \rho_\delta - g(\rho_\delta)) \operatorname{div} u_\delta dy] dx dt \\
& \quad + \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[\operatorname{div} (g(\rho_\delta) u_\delta)] dx dt \\
& \quad - \int_0^T \psi \int_{\mathcal{O}} (\nabla Q_\delta \otimes \nabla Q_\delta - F(Q_\delta) \mathbf{I}_3) : \nabla \mathcal{B}[g(\rho_\delta) - \oint_{\mathcal{O}} g(\rho_\delta) dy] dx dt \\
& \quad + \sigma_* \int_0^T \psi \int_{\mathcal{O}} c_\delta^2 Q_\delta : \nabla \mathcal{B}[g(\rho_\delta) - \oint_{\mathcal{O}} g(\rho_\delta) dy] dx dt \\
& \quad - \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (Q_\delta \Delta Q_\delta - \Delta Q_\delta Q_\delta) \cdot \mathcal{B}[g(\rho_\delta) - \oint_{\mathcal{O}} g(\rho_\delta) dy] dx dt \\
& = \sum_{j=1}^{10} \mathcal{J}_j.
\end{aligned} \tag{5.8}$$

For example, we have the following estimates: Since $\rho_\delta \in L^{\beta+1}(\mathcal{O}_T)$ and $g'(z) = 0$ for sufficiently large z , we have

$$\begin{aligned}
& \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\delta^\beta (S_m[g(\rho_\delta)] - g(\rho_\delta)) dx dt \right| \\
& \leq C \|\rho_\delta\|_{L_{t,x}^{\beta+1}}^\beta \|S_m[g(\rho_\delta)] - g(\rho_\delta)\|_{L_{t,x}^{\beta+1}}
\end{aligned}$$

$$\leq C(\delta) \|S_m[g(\rho_\delta)] - g(\rho_\delta)\|_{L_{t,x}^{\beta+1}} \rightarrow 0 \quad \text{as } m \rightarrow \infty.$$

By the property of \mathcal{B} in (4.2), we have

$$\begin{aligned} \|\partial_j B_i [S_m[g(\rho_\delta)] - \int_{\mathcal{O}} S_m[g(\rho_\delta)] dy]\|_{L_t^\infty L_x^3} &\leq C \|S_m[g(\rho_\delta)] - \int_{\mathcal{O}} S_m[g(\rho_\delta)] dx\|_{L_t^\infty L_x^3} \\ &\leq C \|g(\rho_\delta)\|_{L_t^\infty L_x^3} \leq C(\delta), \end{aligned}$$

where $C(\delta)$ is independent of m . This shows that

$$\partial_j B_i [S_m[g(\rho_\delta)] - \int_{\mathcal{O}} S_m[g(\rho_\delta)] dy] \xrightarrow{*} G \quad \text{weak-star in } L_t^\infty L_x^3.$$

On the other hand, since

$$\begin{aligned} &\|\partial_j B_i [S_m[g(\rho_\delta)] - g(\rho_\delta) - \int_{\mathcal{O}} (S_m[g(\rho_\delta)] - g(\rho_\delta)) dx]\|_{L_t^2 L_x^2} \\ &\leq C \|S_m[g(\rho_\delta)] - g(\rho_\delta) - \int_{\mathcal{O}} (S_m[g(\rho_\delta)] - g(\rho_\delta)) dx\|_{L_t^2 L_x^2} \\ &\leq C \|S_m[g(\rho_\delta)] - g(\rho_\delta)\|_{L_t^2 L_x^2} \rightarrow 0, \end{aligned}$$

we have

$$\partial_j B_i [S_m[g(\rho_\delta)] - \int_{\mathcal{O}} S_m[g(\rho_\delta)] dy] \xrightarrow{*} \partial_j B_i [g(\rho_\delta) - \int_{\mathcal{O}} g(\rho_\delta) dy] \quad \text{weak-star in } L_t^\infty L_x^3.$$

Then

$$\begin{aligned} \mathcal{I}_5 &= \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta^i u_\delta^j \partial_j B_i [S_m[g(\rho_\delta)] - \int_{\mathcal{O}} S_m[g(\rho_\delta)] dy] dx dt \\ &\rightarrow \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta^i u_\delta^j \partial_j B_i [g(\rho_\delta) - \int_{\mathcal{O}} g(\rho_\delta) dy] dx dt = \mathcal{J}_5. \end{aligned}$$

By (4.3), we have

$$\begin{aligned} \mathcal{B}[\operatorname{div}(S_m[g(\rho_\delta)]u_\delta)] &\rightarrow \mathcal{B}[\operatorname{div}(g(\rho_\delta)u_\delta)] \quad \text{strongly in } L_t^{\frac{3}{2}} L_x^2, \\ \mathcal{B}[\operatorname{div}(S_m[g(\rho_\delta)]u_\delta)] &\rightharpoonup G \quad \text{weakly in } L^2(Q_T). \end{aligned}$$

Thus, we obtain

$$\mathcal{B}[\operatorname{div}(S_m[g(\rho_\delta)]u_\delta)] \rightharpoonup \mathcal{B}[\operatorname{div}(g(\rho_\delta)u_\delta)] \quad \text{weakly in } L^2(Q_T).$$

Moreover, since $\psi \rho_\delta u_\delta \in L^2(Q_T)$, we have

$$\mathcal{I}_8 = \int_0^T \int_{\mathcal{O}} \psi \rho_\delta u_\delta \cdot \mathcal{B}[\operatorname{div}(S_m[g(\rho_\delta)]u_\delta)] dx dt \rightarrow \int_0^T \int_{\mathcal{O}} \psi \rho_\delta u_\delta \cdot \mathcal{B}[\operatorname{div}(g(\rho_\delta)u_\delta)] dx dt = \mathcal{J}_7$$

as $m \rightarrow \infty$. Then, as we mentioned before, we can use a sequence of functions $\{z^\theta \chi_n(z)\}$ to approximate $g(z) = z^\theta$ to obtain

$$\int_0^T \int_{\mathcal{O}} \psi (\rho_\delta^{\gamma+\theta} + \delta \rho_\delta^{\beta+\theta}) dx dt = \sum_{i=1}^{10} \mathcal{J}_i, \quad (5.9)$$

where $g(z)$ is substituted by z^θ in every \mathcal{J}_i , $i = 1, 2, \dots, 10$.

Next, we estimate the terms on the right-hand side of (5.9) to obtain the condition for θ , for which the universal constant C is independent of δ :

$$\begin{aligned}
|\mathcal{J}_1| &= \left| \int_0^T \psi \int_{\mathcal{O}} (\rho_\delta^\gamma + \delta \rho_\delta^\beta) dx \int_{\mathcal{O}} \rho_\delta^\theta dx dt \right| \\
&\leq C \|\psi\|_{L^\infty(0,T)} (\|\rho_\delta\|_{L_t^\infty L_x^\gamma}^\gamma + \delta \|\rho_\delta\|_{L_t^\infty L_x^\beta}^\beta) \|\rho_\delta\|_{L_t^\infty L_x^\theta}^\theta \leq C \|\psi\|_{L^\infty(0,T)}, \\
|\mathcal{J}_2| &= (\mu + \nu) \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\delta^\theta \operatorname{div} u_\delta dx dt \right| \leq C \|\nabla u_\delta\|_{L^2(Q_T)} \|\rho_\delta^\theta\|_{L_t^\infty L_x^2} \\
&\leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{\gamma}{2}, \\
|\mathcal{J}_3| &= \left| \int_0^T \psi_t \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx dt \right| \\
&\leq C \int_0^T |\psi_t| \|\sqrt{\rho_\delta}\|_{L_x^{2\gamma}} \|\sqrt{\rho_\delta} u_\delta\|_{L_x^2} \|\mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^{\frac{2\gamma}{\gamma-1}}} dt \\
&\leq C \int_0^T |\psi_t| \|\rho_\delta\|_{L_x^\gamma}^{\frac{1}{2}} \|\sqrt{\rho_\delta} u_\delta\|_{L_x^2} \|\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy\|_{L_x^{\frac{2\gamma}{\gamma-1}}} dt \\
&\leq CT \|\rho_\delta\|_{L_t^\infty L_x^\gamma}^{\frac{1}{2}} \|\sqrt{\rho_\delta} u_\delta\|_{L_t^\infty L_x^2} \|\rho_\delta\|_{L_t^\infty L_x^{\frac{2\gamma}{\gamma-1}\theta}}^\theta \|\psi\|_{L^1(0,T)} \\
&\leq C(T) \|\psi\|_{L^1(0,T)} \quad \text{if } \theta \leq \frac{\gamma-1}{2}, \\
|\mathcal{J}_4| &= \mu \left| \int_0^T \psi \int_{\mathcal{O}} \partial_j u_\delta^i \partial_j B_i [\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx dt \right| \\
&\leq \mu \int_0^T |\psi| \|\nabla u_\delta\|_{L_x^2} \|\nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^2} dt \\
&\leq \mu \int_0^T |\psi| \|\nabla u_\delta\|_{L_x^2} \|\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy\|_{L_x^2} dt \leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{\gamma}{2}, \\
|\mathcal{J}_5| &= \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta^i u_\delta^j \partial_j B_i [\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx dt \right| \\
&\leq C \int_0^T |\psi| \|\rho_\delta\|_{L_x^\gamma} \|u_\delta\|_{L_x^6}^2 \|\nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^{\frac{3\gamma}{2\gamma-3}}} dt \\
&\leq C \int_0^T |\psi| \|\rho_\delta\|_{L_x^\gamma} \|u_\delta\|_{L_x^6}^2 \|\rho_\delta^\theta\|_{L_x^{\frac{3\gamma}{2\gamma-3}}} dt \leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{2\gamma}{3} - 1, \\
|\mathcal{J}_6| &= (1 - \theta) \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[\rho_\delta^\theta \operatorname{div} u_\delta - \int_{\mathcal{O}} \rho_\delta^\theta \operatorname{div} u_\delta dy] dx dt \right| \\
&\leq C \int_0^T |\psi| \|\rho_\delta\|_{L_x^\gamma} \|u_\delta\|_{L_x^6} \|\mathcal{B}[\rho_\delta^\theta \operatorname{div} u_\delta - \int_{\mathcal{O}} \rho_\delta^\theta \operatorname{div} u_\delta dy]\|_{L_x^{\frac{6\gamma}{5\gamma-6}}} dt \\
&\leq C \int_0^T |\psi| \|u_\delta\|_{L_x^6} \|\mathcal{B}[\rho_\delta^\theta \operatorname{div} u_\delta - \int_{\mathcal{O}} \rho_\delta^\theta \operatorname{div} u_\delta dy]\|_{W_x^{1,p}} dt
\end{aligned}$$

$$\begin{aligned}
&\leq C \int_0^T |\psi| \|u_\delta\|_{L_x^6} \|\rho_\delta^\theta \operatorname{div} u_\delta\|_{L_x^p} dt \\
&\leq C \int_0^T |\psi| \|u_\delta\|_{L_x^6} \|\nabla u_\delta\|_{L_x^2} \|\rho_\delta^\theta\|_{L_x^q} dt \\
&\leq C \int_0^T |\psi| \|u_\delta\|_{H_x^1}^2 \|\rho_\delta\|_{L_x^\gamma} dt \leq C \|\psi\|_{L^\infty(0,T)},
\end{aligned}$$

where $p = \frac{6\gamma}{7\gamma-6}$ if $\gamma < 6$, $p = \frac{3}{2}$ if $\gamma \geq 6$; $q = \frac{3\gamma}{2\gamma-3}$ if $\gamma < 6$, $q = 6$ if $\gamma \geq 6$; $\theta \leq \frac{2}{3}\gamma - 1$ and $\theta \leq 1$.

$$\begin{aligned}
|\mathcal{J}_7| &= \left| \int_0^T \psi \int_{\mathcal{O}} \rho_\delta u_\delta \cdot \mathcal{B}[\operatorname{div}(\rho_\delta^\theta u_\delta)] dx dt \right| \leq \int_0^T |\psi| \|\rho_\delta\|_{L_x^\gamma} \|u_\delta\|_{L_x^6} \|\mathcal{B}[\operatorname{div}(\rho_\delta^\theta u_\delta)]\|_{L_x^{\frac{6\gamma}{5\gamma-6}}} dt \\
&\leq \int_0^T |\psi| \|u_\delta\|_{L_x^6} \|\rho_\delta^\theta u_\delta\|_{L_x^{\frac{6\gamma}{5\gamma-6}}} dt \leq \int_0^T |\psi| \|u_\delta\|_{L_x^6}^2 \|\rho_\delta^\theta\|_{L_x^{\frac{3\gamma}{2\gamma-3}}} dt \\
&\leq C \|u_\delta\|_{L_t^2 H_x^1}^2 \|\rho_\delta\|_{L_t^\infty L_x^\gamma}^\theta \|\psi\|_{L^\infty(0,T)} \\
&\leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{2\gamma-3}{3}.
\end{aligned}$$

Similarly to (4.6), we have

$$\begin{aligned}
&\|B[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^\infty} \\
&\leq C_1 \|\nabla B[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^\gamma}^{\frac{3}{\gamma}} \|B[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^\gamma}^{1-\frac{3}{\gamma}} + C_2 \|B[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^\gamma} \\
&\leq C \|\rho_\delta^\theta\|_{L_x^\gamma}.
\end{aligned}$$

Then we have

$$\begin{aligned}
|\mathcal{J}_8| &= \left| \int_0^T \psi \int_{\mathcal{O}} (\nabla Q_\delta \otimes \nabla Q_\delta - F(Q_\delta) I_3) : \nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx dt \right| \\
&\leq C \int_0^T |\psi| \left(\|\nabla Q_\delta\|_{L_x^{\frac{10}{3}}}^2 \|\nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^{\frac{5}{2}}} + \left| \int_{\mathcal{O}} F(Q_\delta) \operatorname{div} \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx \right| \right) dt \\
&\leq C \int_0^T |\psi| \left(\|\nabla Q_\delta\|_{L_x^{\frac{10}{3}}}^2 \|\nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^{\frac{5}{2}}} + \left| \int_{\mathcal{O}} F(Q_\delta) (\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy) dx \right| \right) dt \\
&\leq C \int_0^T |\psi| \left(\|\nabla Q_\delta\|_{L_x^{\frac{10}{3}}}^2 \|\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy\|_{L_x^{\frac{5}{2}}} + (\|Q_\delta\|_{L_x^5}^2 + \|Q_\delta\|_{L_x^{10}}^4) \|\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy\|_{L_x^{\frac{5}{3}}} \right) dt \\
&\leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{2}{5},
\end{aligned}$$

$$\begin{aligned}
|\mathcal{I}_9| &= \sigma_* \left| \int_0^T \psi \int_{\mathcal{O}} c_\delta^2 Q_\delta \cdot \nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx dt \right| \\
&\leq C \|c_\delta\|_{L_{t,x}^\infty}^2 \int_0^T |\psi| \|Q_\delta\|_{L_x^2} \|\nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^2} dt \\
&\leq C \int_0^T |\psi| \|Q_\delta\|_{L_x^2} \|\rho_\delta^\theta\|_{L_x^2} dt \leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{\gamma}{2},
\end{aligned}$$

$$\begin{aligned}
|\mathcal{I}_{10}| &= \left| \int_0^T \psi \int_{\mathcal{O}} \nabla \cdot (Q_\delta \Delta Q_\delta - \Delta Q_\delta Q_\delta) \cdot \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy] dx dt \right| \\
&\leq \int_0^T |\psi| \|Q_\delta\|_{L_x^4} \|\Delta Q_\delta\|_{L_x^2} \|\nabla \mathcal{B}[\rho_\delta^\theta - \int_{\mathcal{O}} \rho_\delta^\theta dy]\|_{L_x^4} dt \\
&\leq \int_0^T |\psi| \|Q_\delta\|_{L_x^4} \|\Delta Q_\delta\|_{L_x^2} \|\rho_\delta^\theta\|_{L_x^4} dt \leq C \|\psi\|_{L^\infty(0,T)} \quad \text{if } \theta \leq \frac{1}{4}.
\end{aligned}$$

Combining the above estimates with (5.9), we have

$$\int_0^T \int_{\mathcal{O}} \psi (\rho_\delta^{\gamma+\theta} + \delta \rho_\delta^{\beta+\theta}) dx dt \leq C (\|\psi\|_{L^\infty(0,T)} + \|\psi_t\|_{L^1(0,T)}). \quad (5.10)$$

Then we can again take $\psi = \psi^k \in \mathcal{D}(0, T)$ in (5.10), with $\|\psi^k\|_{L^\infty(0,T)} + \|\psi_t^k\|_{L^1(0,T)} \leq C$ uniformly bounded independent of k , and $\psi^k(t) \rightarrow 1$ for $t \in (0, T)$ as $k \rightarrow \infty$, and then let $k \rightarrow \infty$ in the resulting (5.10) to conclude our desired result. \square

5.2. The limit passage and the effective viscous flux. We infer from the uniform estimates (4.38)–(4.45) in Proposition 4.1 and Lemma 5.1 that, as $\delta \rightarrow 0$,

$$\begin{aligned}
c_\delta &\rightharpoonup c && \text{in } L^2(0, T; H^1(\mathcal{O})), \\
\rho_\delta &\rightarrow \rho && \text{in } C([0, T]; L_{\text{weak}}^\gamma(\mathcal{O})), \\
u_\delta &\rightharpoonup u && \text{in } L^2(0, T; H_0^1(\mathcal{O})), \\
\rho_\delta u_\delta &\rightarrow \rho u && \text{in } C([0, T]; L_{\text{weak}}^{\frac{2\gamma}{\gamma+1}}(\mathcal{O})), \\
Q_\delta &\rightharpoonup Q && \text{in } L^2(0, T; H^2(\mathcal{O})), \\
Q_\delta &\rightarrow Q && \text{in } L^2(0, T; H^1(\mathcal{O})), \\
\rho_\delta^\gamma &\rightarrow \overline{\rho^\gamma} && \text{in } L^{\frac{\gamma+\theta}{\gamma}}(\mathcal{O}_T).
\end{aligned}$$

Moreover, we have

$$\begin{aligned}
\rho_\delta u_\delta \otimes u_\delta &\rightarrow \rho u \otimes u && \text{in } \mathcal{D}'(\mathcal{O}_T), \\
F(Q_\delta) \mathbf{I}_3 - \nabla Q_\delta \odot \nabla Q_\delta + (Q_\delta \Delta Q_\delta - \Delta Q_\delta Q_\delta) + \sigma_* c_\delta^2 Q_\delta \\
&\rightarrow F(Q) \mathbf{I}_3 - \nabla Q \odot \nabla Q + (Q \Delta Q - \Delta Q Q) + \sigma_* c^2 Q, \\
\delta \rho_\delta^\beta &\rightarrow 0 && \text{in } L^1(\mathcal{O}_T).
\end{aligned}$$

Then the limit functions (c, ρ, u, Q) satisfy

$$c_t + u \cdot \nabla c = D_0 \Delta c, \quad (5.11)$$

$$\partial_t \rho + \nabla \cdot (\rho u) = 0, \quad (5.12)$$

$$(\rho u)_t + \nabla \cdot (\rho u \otimes u) + \nabla \overline{\rho^\gamma} = \mu \Delta u + (\nu + \mu) \nabla \operatorname{div} u + \nabla \cdot (Q \Delta Q - \Delta Q Q) \quad (5.13)$$

$$+ \nabla \cdot (F(Q) \mathbf{I}_3 - \nabla Q \odot \nabla Q) + \sigma_* \nabla \cdot (c^2 Q), \quad (5.14)$$

$$\partial_t Q + (u \cdot \nabla) Q + Q \Omega - \Omega Q = \Gamma H[Q, c], \quad (5.15)$$

in $\mathcal{D}'(\mathcal{O}_T)$, with the initial-boundary conditions (1.6)–(1.8), due to (5.2) and (5.4).

Next, in order to prove that (c, ρ, u, Q) is a weak solution of (1.2)–(1.8), we only need to show that $\bar{\rho}^\gamma = \rho^\gamma$ *a.e.* in \mathcal{O}_T , or equivalently, the strong convergence of ρ_δ in $L^1(\mathcal{O})$. As in §4, we need to show that (ρ, u) is a renormalized solution of (5.12). From Lemma 5.1, the best estimate is $\rho \in L^{\gamma + \frac{2}{3}\gamma - 1}(\mathcal{O}_T)$. Then $\gamma > \frac{3}{2}$ is not enough to guarantee that ρ is square integrable, and that (ρ, u) is a renormalized solution by Lemma A.4 in Appendix A. In order to deal with this difficulty, we introduce the cut-off function $T_k(z) = kT(\frac{z}{k})$ for $z \in \mathbb{R}$, $k = 1, 2, \dots$, where T is a smooth and concave function satisfying

$$T(z) = \begin{cases} z & \text{if } z \leq 1, \\ 2 & \text{if } z \geq 3. \end{cases}$$

Since (ρ_δ, u_δ) is a renormalized solution of (4.35), taking $g(z) = T_k(z)$, we obtain

$$\partial_t(T_k(\rho_\delta)) + \operatorname{div}(T_k(\rho_\delta)u_\delta) + (T'_k(\rho_\delta) - T_k(\rho_\delta))\operatorname{div} u_\delta = 0 \quad \text{in } \mathcal{D}'((0, T) \times \mathbb{R}^3), \quad (5.16)$$

where $T_k(\rho_\delta) \in L^\infty(\mathcal{O}_T)$ for any fixed k . Then

$$T_k(\rho_\delta) \xrightarrow{*} \overline{T_k(\rho)} \quad \text{in } L^\infty(\mathcal{O}_T).$$

Moreover, since $\partial_t T_k(\rho_\delta)$ satisfies (5.16), as before, we have

$$T_k(\rho_\delta) \rightarrow \overline{T_k(\rho)} \quad \text{in } C([0, T]; L^p_{\text{weak}}(\mathcal{O})), \quad \forall 1 \leq p < \infty.$$

Letting $\delta \rightarrow 0$, it yields

$$\partial_t \overline{T_k(\rho)} + \operatorname{div}(\overline{T_k(\rho)}u) + \overline{(T'_k(\rho) - T_k(\rho))\operatorname{div} u} = 0 \quad \text{in } \mathcal{D}'(\mathcal{O}_T),$$

where

$$(T'_k(\rho_\delta) - T_k(\rho_\delta))\operatorname{div} u_\delta \rightarrow \overline{(T'_k(\rho) - T_k(\rho))\operatorname{div} u} \quad \text{in } L^2(\mathcal{O}_T).$$

5.3. The effective viscous flux. Similarly to Lemma 4.2, we define the effective viscous flux as $\tilde{\mathfrak{E}}_\delta := \rho_\delta^\gamma - (\nu + 2\mu)\operatorname{div} u_\delta$, and its correspondingly weak convergence limit $\tilde{\mathfrak{E}} := \bar{\rho}^\gamma - (\nu + 2\mu)\operatorname{div} u$. Then we have

Lemma 5.2. *Assume (ρ_δ, u_δ) is a family of the approximate solutions constructed in Proposition 4.1. Then*

$$\lim_{\delta \rightarrow 0^+} \int_0^T \psi \left(\int_{\mathcal{O}} \phi \tilde{\mathfrak{E}}_\delta T_k(\rho_\delta) dx \right) dt = \int_0^T \psi \left(\int_{\mathcal{O}} \phi \tilde{\mathfrak{E}} \overline{T_k(\rho)} dx \right) dt,$$

for any $\psi \in \mathcal{D}(0, T)$ and $\phi \in \mathcal{D}(\mathcal{O})$.

5.4. Renormalized solutions. The following lemma implies that $T_k(\rho) - \overline{T_k(\rho)} \in L^2(\mathcal{O}_T)$, which helps us establish that the limit function (ρ, u) is a renormalized solution.

Lemma 5.3 (The amplitude of oscillations). *There exists a constant C , independent of k , such that*

$$\limsup_{\delta \rightarrow 0} \|T_k(\rho_\delta) - T_k(\rho)\|_{L^{\gamma+1}(\mathcal{O}_T)} \leq C. \quad (5.17)$$

The proof of this lemma is the same as that for Lemma 4.3 of Subsection 4.4 in [15].

Remark 5.1. By the concavity of the norm, we know from Lemma 5.3 that

$$\|\overline{T_k(\rho)} - T_k(\rho)\|_{L^{\gamma+1}(\mathcal{O}_T)} \leq C. \quad (5.18)$$

From the proof of Lemma 5.3, we have

$$0 \leq \limsup_{\delta \rightarrow 0} \|T_k(\rho_\delta) - T_k(\rho)\|_{L^{\gamma+1}(\mathcal{O}_T)}^{\gamma+1} \leq \lim_{\delta \rightarrow 0} \int_0^T \int_{\mathcal{O}_T} (\rho_\delta^\gamma T_k(\rho_\delta) - \overline{\rho^\gamma} \overline{T_k(\rho)}) dx dt. \quad (5.19)$$

Based on the uniform estimate of the amplitude of oscillations in Lemma 5.3, we see that the limit function (ρ, u) satisfies (5.12) in the renormalized sense.

Lemma 5.4. *The limit function (ρ, u) is a renormalized solution to (5.12), that is,*

$$\partial_t g(\rho) + \operatorname{div}(g(\rho)u) + (g'(\rho)\rho - g(\rho))\operatorname{div} u = 0 \quad \text{in } \mathcal{D}'((0, T) \times \mathbb{R}^3), \quad (5.20)$$

for any $g \in C^1(\mathbb{R})$ with the property $g'(z) \equiv 0$ when $z \geq M$ for sufficiently large constant M , provided (ρ, u) are prolonged zero outside \mathcal{O} .

The detailed proof can be found in Subsection 4.5 of [15] for Lemma 4.4 there.

5.5. Strong convergence of density ρ_δ . We now give an outline of the proof for the strong convergence of density ρ_δ , i.e., $\overline{\rho^\gamma} = \rho^\gamma$, where $\rho_\delta^\gamma \rightarrow \overline{\rho^\gamma}$ in $L^{\frac{\gamma+\theta}{\gamma}}(\mathcal{O}_T)$ for the completeness of the proof of Theorem 1.1.

Introduce a family of function in $C^1(\mathbb{R}^+) \cap C[0, \infty)$:

$$L_k(z) = \begin{cases} z \log z & \text{for } 0 \leq z < k, \\ z \log k + z \int_k^z \frac{T_k(s)}{s^2} ds & \text{for } z \geq k. \end{cases}$$

By the construction of L_k , we know that L_k is a linear function for large z . In particular, we see that, for $z \geq 3k$,

$$L_k(z) = \beta_k z - 2k$$

with

$$\beta_k = \log k + \int_k^{3k} \frac{T_k(s)}{s^2} ds + \frac{2}{3}.$$

Then, if $g_k(z) := L_k(z) - \beta_k z$, we obtain that $g_k(z) \in C^1(\mathbb{R}^+) \cap C[0, \infty)$, $g'_k(z) = 0$ for z is sufficiently large, and

$$g'_k(z)z - g_k(z) = T_k(z).$$

By Proposition 4.1 and Lemma 5.4, we know that (ρ_δ, u_δ) and (ρ, u) are renormalized solutions to (5.12). Then we substitute function g by g_k in the definition of renormalized solutions and take the difference of these two equations to obtain

$$\partial_t (L_k(\rho_\delta) - L_k(\rho)) + \operatorname{div}(L_k(\rho_\delta)u_\delta - L_k(\rho)u) + T_k(\rho_\delta)\operatorname{div} u_\delta - T_k(\rho)\operatorname{div} u = 0 \quad (5.21)$$

in $\mathcal{D}'((0, T) \times \mathbb{R}^3)$.

Since L_k is linear when z is large, $L_k(\rho_\delta)$ is uniformly bounded with respect to δ in $L^\infty(0, T; L^\gamma(\mathcal{O}))$ so that

$$L_k(\rho_\delta) \xrightarrow{*} \overline{L_k(\rho)} \quad \text{in } L_t^\infty L_x^\gamma \text{ as } \delta \rightarrow 0.$$

Moreover, since $L_k(\rho_\delta)$ is a renormalized solution, similarly as before, we have

$$L_k(\rho_\delta) \rightarrow \overline{L_k(\rho)} \quad \text{in } C([0, T], L_{\text{weak}}^\gamma(\mathcal{O})) \cap C([0, T]; H^{-1}(\mathcal{O})) \text{ as } \delta \rightarrow 0. \quad (5.22)$$

Now, using function $\phi(x) \in \mathcal{D}(\mathcal{O})$ to test (5.21) and then integrate over $(0, t)$, we have

$$\begin{aligned} & \int_{\mathcal{O}} (L_k(\rho_\delta) - L_k(\rho))(t, x) \phi(x) dx \\ &= \int_{\mathcal{O}} (L_k(\rho_{0,\delta}) - L_k(\rho_0)) \phi dx + \int_0^t \int_{\mathcal{O}} (L_k(\rho_\delta) u_\delta - L_k(\rho) u) \cdot \nabla \phi dx ds \\ & \quad - \int_0^t \int_{\mathcal{O}} (T_k(\rho_\delta) \operatorname{div} u_\delta - T_k(\rho) \operatorname{div} u) \phi dx ds. \end{aligned}$$

Sending $\delta \rightarrow 0$, we have

$$\begin{aligned} & \int_{\mathcal{O}} (\overline{L_k(\rho)} - L_k(\rho))(t, x) \phi(x) dx \\ &= \int_0^t \int_{\mathcal{O}} (\overline{L_k(\rho)} - L_k(\rho)) u \cdot \nabla \phi dx ds - \lim_{\delta \rightarrow 0} \int_0^t \int_{\mathcal{O}} (T_k(\rho_\delta) \operatorname{div} u_\delta - T_k(\rho) \operatorname{div} u) \phi dx ds. \end{aligned}$$

Taking $\phi = \phi_m$ with $\phi_m(z) \rightarrow 1_{\mathcal{O}}(z)$ in the above equation and sending $m \rightarrow \infty$, we have

$$\int_{\mathcal{O}} (\overline{L_k(\rho)} - L_k(\rho))(t) dx = \int_0^t \int_{\mathcal{O}} T_k(\rho) \operatorname{div} u dx ds - \lim_{\delta \rightarrow 0} \int_0^t \int_{\mathcal{O}} T_k(\rho_\delta) \operatorname{div} u_\delta dx ds. \quad (5.23)$$

Then, using Lemmas 5.2–5.3, we find that the right-hand side of the above equation is non-positive, which yields

$$\lim_{k \rightarrow \infty} \int_{\mathcal{O}} (\overline{L_k(\rho)} - L_k(\rho))(t) dx \leq 0 \quad \text{for } t \in [0, T]. \quad (5.24)$$

Moreover, by the definition of L_k and the absolute continuity of $\rho \log \rho \in L^1(\mathcal{O}_T)$, we have

$$\begin{aligned} \|L_k(\rho) - \rho \log \rho\|_{L^1(\mathcal{O}_T)} &\leq \int \int_{\{\rho \geq k\}} |L_k(\rho) - \rho \log \rho| dx dt \\ &\leq C \int \int_{\{\rho \geq k\}} |\rho \log \rho| dx dt \rightarrow 0 \quad \text{as } k \rightarrow \infty. \end{aligned} \quad (5.25)$$

Similarly, we have

$$\begin{aligned} \|L_k(\rho_\delta) - \rho_\delta \log \rho_\delta\|_{L^1(\mathcal{O}_T)} &\leq \int \int_{\{\rho_\delta \geq k\}} |L_k(\rho_\delta) - \rho_\delta \log \rho_\delta| dx dt \\ &\leq \int \int_{\{\rho_\delta \geq k\}} \frac{\log k + \int_k^{\rho_\delta} \frac{T_k(s)}{s^2} ds + \log \rho_\delta}{\rho_\delta^{\gamma-1}} \rho_\delta^\gamma dx dt \\ &\leq C(\varepsilon) k^{1+\varepsilon-\gamma} \int \int_{\{\rho_\delta \geq k\}} \rho_\delta^\gamma dx dt \rightarrow 0 \quad \text{as } k \rightarrow \infty. \end{aligned}$$

This, together with the lower-semicontinuity of the norm, we have

$$\|\overline{L_k(\rho)} - \overline{\rho \log \rho}\|_{L^1(\mathcal{O}_T)} \leq \liminf_{\delta \rightarrow 0} \|L_k(\rho_\delta) - \rho_\delta \log \rho_\delta\|_{L^1(\mathcal{O}_T)} \rightarrow 0 \quad \text{as } k \rightarrow \infty. \quad (5.26)$$

Finally, combining (5.24)–(5.26), we have

$$\int_{\mathcal{O}} (\overline{\rho \log \rho} - \rho \log \rho)(t) dx \leq 0.$$

Moreover, since $\overline{\rho \log \rho} \geq \rho \log \rho$, we see that $\overline{\rho \log \rho} = \rho \log \rho$ for *a.e.* $(t, x) \in \mathcal{O}_T$. In addition, by the restrict concavity of function $z \log z$, we have

$$\rho_\delta \rightarrow \rho \quad \text{in } L^p(\mathcal{O}_T) \text{ for any } p \in [1, \gamma + \theta).$$

Then we conclude

$$\overline{\rho^\gamma} = \rho^\gamma, \quad \text{a.e.}$$

Therefore, we complete the proof of Theorem 1.1.

APPENDIX A. PRELIMINARIES

In this appendix, we collect some important results that we use extensively in this paper.

In order to deal with the highest derivatives of u in (1.4) (and the corresponding ones in the approximation systems) and the highest derivatives of Q in (1.5) (and the corresponding ones in the approximation systems), we need the following lemma whose proof can be found in the proof of Lemma A.1 in [9].

Lemma A.1. *Let Q and Q' be two 3×3 symmetric matrices, and let $\Omega = \frac{1}{2}(\nabla u - (\nabla u)^\top)$ be the vorticity with $(\nabla u)_{\alpha\beta} = \partial_\beta u_\alpha$. Then*

$$(\Omega Q' - Q' \Omega, \Delta Q) - (\nabla \cdot (Q' \Delta Q - \Delta Q Q'), u) = 0.$$

Lemma A.2 (Aubin-Lions lemma [1]). *Let X_0, X , and X_1 be three Banach spaces with $X_0 \subseteq X \subseteq X_1$, X_0 compactly embedded in X , and X continuously embedded in X_1 . For $1 \leq p, q \leq \infty$, let*

$$W = \{u \in L^p(0, T; X_0) : \dot{u} \in L^q(0, T; X_1)\}.$$

Then

- (i) *If $p < \infty$, then the embedding of W into $L^p(0, T; X)$ is compact;*
- (ii) *If $p = \infty$ and $q > 1$, then the embedding of W into $C(0, T; X)$ is compact.*

Lemma A.3 (Gagliardo-Nirenberg interpolation inequality [41]). *Let $1 \leq q, r \leq \infty$ and $0 \leq j < m$. Then the following inequality holds:*

$$\|D^j u\|_{L^p} \leq C_1 \|D^m u\|_{L^r}^a \|u\|_{L^q}^{1-a} + C_2 \|u\|_{L^s}$$

for any function $u : \mathcal{O} \rightarrow \mathbb{R}$ defined on a bounded Lipschitz domain $\mathcal{O} \subseteq \mathbb{R}^3$, where

$$\frac{1}{p} = \frac{j}{3} + a\left(\frac{1}{r} - \frac{m}{3}\right) + (1-a)\frac{1}{q}, \quad \frac{j}{m} \leq a \leq 1,$$

$s > 0$ is arbitrary, and C_1 and C_2 depend only on \mathcal{O} and m .

Next, we introduce a sufficient condition for a solution (ρ, u) to be a renormalized solution.

Lemma A.4. *Let $\mathcal{O} \subseteq \mathbb{R}^3$ be a bounded domain, and let $\rho \in L^2(\mathcal{O}_T)$ and $u \in L^2(0, T; H^1(\mathcal{O}))$ such that*

$$\partial_t \rho + \nabla \cdot (\rho u) = 0 \quad \text{in } \mathcal{D}'(\mathcal{O}_T).$$

Then

$$\partial_t g(\rho) + \operatorname{div}(g(\rho)u) + (g'(\rho)\rho - g(\rho))\operatorname{div} u = 0 \quad \text{in } \mathcal{D}'(\mathcal{O}_T) \quad (\text{A.1})$$

for any $g \in C^1(\mathbb{R})$ with the property: $g'(z) \equiv 0$, when $z \geq M$ for sufficiently large constant M , i.e., (ρ, u) is a renormalized solution.

Definition A.1. The metric space $C([0, T]; X_{\text{weak}}^*)$ contains all the functions $v : [0, T] \mapsto X^*$ which are continuous with respect to the weak topology. We say

$$v_n \rightarrow v \quad \text{in } C([0, T]; X_{\text{weak}}^*),$$

if $(v_n(t), \phi) \rightarrow (v(t), \phi)$ uniformly with respect to $t \in [0, T]$ for any $\phi \in X$.

In the following corollary, we introduce a sufficient condition for a family of functions to converge in $C([0, T]; X_{\text{weak}}^*)$ (see Corollary 2.1 in [16]).

Corollary A.1. *Let X be a separable Banach space. Assume that $v_n : [0, T] \rightarrow X^*$, $n = 1, 2, \dots$, is a sequence of measurable functions such that*

$$\text{ess sup}_{t \in [0, T]} \|v_n(t)\|_{X^*} \leq M, \quad \text{uniformly in } n = 1, 2, \dots$$

Moreover, let the family of functions:

$$(v_n, \phi) : t \mapsto (v_n(t), \phi), \quad t \in [0, T], n = 1, 2, \dots,$$

be equi-continuous for any fixed ϕ belonging to a dense subset in X .

Then $v_n \in C([0, T]; X_{\text{weak}}^*)$ for any $n = 1, 2, \dots$, and there exist $v \in C([0, T]; X_{\text{weak}}^*)$ and a subsequence (still denoted) v_n such that

$$v_n \rightarrow v \quad \text{in } C([0, T]; X_{\text{weak}}^*) \quad \text{as } n \rightarrow \infty.$$

Proposition A.1. *Let \mathcal{O} be a bounded domain in \mathbb{R}^3 . If $v_n \in L^\infty(0, T; L^p(\mathcal{O}))$ for $p > \frac{6}{5}$, and $v_n \rightarrow v$ in $C([0, T], L_{\text{weak}}^p(\mathcal{O}))$, then $v_n \rightarrow v$ in $C([0, T], H^{-1}(\mathcal{O}))$.*

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MATHEMATICAL INSTITUTE, UNIVERSITY OF OXFORD, OXFORD OX2 6GG, UK.

E-mail address: `chengq@maths.ox.ac.uk`

DEPARTMENT OF MATHEMATICAL SCIENCES, UNIVERSITY OF BATH, BATH, BA2 7AY, UK.

E-mail address: `a.majumdar@bath.ac.uk`

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF PITTSBURGH, PITTSBURGH, PA 15260, USA.

E-mail address: `dwang@math.pitt.edu`

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF PITTSBURGH, PITTSBURGH, PA 15260, USA.

E-mail address: `roz14@pitt.edu`